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# Spectroscopic Observations of Some S-Type Symbiotic Stars 

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#### Abstract

Observations of a group of S-type symbiotic stars are presented: Hen 828, Hen 905, Hen 1103, Hen 1213, KX TrA (Cn 1-2), AE Ara (PC 18), AS 255, AS 270, AS 289, Y CrA, AS 304, AS 316, AS 327, and CD $-43^{\circ} 14304$. The observations were made in three different epochs of observation: 1987, 1991, and 1995, except for Y CrA, which was observed in 1990. Several parameters are obtained for each object: spectral types of the late components; distances derived from these spectral types; interstellar reddening as determined from the Balmer decrement, considering optical depth effects; electron temperatures and densities; helium abundances, also determined considering optical depth effects; and characteristics of the hot components. It has been found that Hen 828 and Hen 1103 were possibly having an outburst in 1991 ; the same could be true for AS 270 in 1995 and CD $-43^{\circ} 14304$ in 1991. Another interesting object is Hen 905 , which is surrounded by a nebulosity that could be a ring, reaching a distance of about $2^{\prime}$, with emission mainly in $\mathrm{H} \alpha$ and O vi $\lambda 6825$.


## 1. INTRODUCTION

Symbiotic stars are interacting binary systems with three components: an evolved cool giant clearly seen in the red part of the optical spectrum, a hot radiation source, and an ionized nebula that is conspicuous because of its emission lines. The hot component can be a wind-accreting white dwarf, a hot subdwarf, or a disk-accreting main-sequence star; its physical characteristics are similar to those of the central stars of planetary nebulae, with $T^{*} \approx 10^{4}-10^{5} \mathrm{~K}$ and $\log L^{*}=10^{3}-10^{4} L_{\odot}$.

According to the evidence of dust deduced from the nearinfrared colors, symbiotic stars have been divided (Allen 1982a, 1982b) in two major types: S types, which are characterized by their IR continua, showing only the presence of a cool star; and D and $\mathrm{D}^{\prime}$ types, which are characterized by infrared radiation dominated by dust with temperatures of the order of $800-1000 \mathrm{~K}$ for D type, or half as great for the D' type (Allen 1982a, 1982b). S-type systems contain highdensity nebulae ( $N_{e} \geq 10^{7} \mathrm{~cm}^{-3}$ ) ionized by a hot source with $T^{*}>70,000 \mathrm{~K} . \mathrm{D}$ and $\mathrm{D}^{\prime}$ types have lower density nebulae ( $N_{e} \leq 10^{7} \mathrm{~cm}^{-3}$ ) and hot sources of lower temperature ( $T^{*} \approx 60,000 \mathrm{~K}$ ).

In a previous paper Gutiérrez-Moreno, Moreno, \& Cortés (1995) presented a diagnostic diagram adequate for

[^0]separating planetary nebulae from symbiotic stars, based only on the behavior of the [ O III ] lines in the optical region. Besides, the diagram separates very well the D - and S-type symbiotic objects, showing that D - and $\mathrm{D}^{\prime}$-type systems have densities of the order of $10^{6}-10^{7} \mathrm{~cm}^{-1}$, whereas S-type systems have densities larger than $10^{7}$ $\mathrm{cm}^{-1}$, coming close to the high-density limit. Several symbiotic stars were observed with the aim of obtaining this diagram. We have already presented an analysis of the observations of a group of symbiotic stars of $D$ and $D^{\prime}$ type (Gutiérrez-Moreno \& Moreno 1996); in this paper we present a similar analysis of the observations of some S-type symbiotic stars.

## 2. OBSERVATIONS AND REDUCTIONS

Details of the objects observed are given in Table 1, which lists, by column, (1) and (2) the name or names of the objects observed; (3) and (4) their positions referred to the equinox 2000.0 ; (5)-(8) the magnitude $K$ and the colors $J-H, H-K$, and $K-L$, taken from Munari et al. (1992) except when indicated; and (9) the spectral types of the late components as listed by Whitelock \& Munari (1992).

The observations were made in five different runs; the details concerning the instrumentation used are listed in Table 2, which is self-explanatory. Table 3 indicates the observing runs corresponding to each object.

For all the observations the slit was oriented in the eastwest direction. All the objects were observed through a

TABLE 1
Observed Objects

| Name <br> (1) | Other <br> (2) | $\begin{aligned} & \text { R.A. } \\ & (2000.0) \\ & (3) \end{aligned}$ | $\begin{gathered} \text { Decl. } \\ (2000.0) \\ (4) \end{gathered}$ | $\begin{gathered} K \\ (5) \end{gathered}$ | $J-H$ <br> (6) | $H-K$ <br> (7) | $K-L$ <br> (8) | Spectral Type <br> (9) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Hen 828 |  | 125058.0 | -575046 | 7.17 | 1.10 | 0.38 | $\ldots$ | M6 |
| Hen 905 |  | 133037.2 | -575817 | 8.39 | 1.12 | 0.35 | $\ldots$ | K4: |
| Hen 1103. |  | 154828.4 | -441900 | 8.40 | 1.00 | 0.30 | $\ldots$ | M0 |
| Hen 1213 ${ }^{\text {a }}$ |  | 163515.2 | -514224 | 6.72 | 1.08 | 0.29 | $\ldots$ | K4 |
| KX $\operatorname{Tr} \mathrm{A} \ldots \ldots \ldots \ldots .$. | Cn 1-2, Hen 1242 | 164435.5 | -62 3713 | 5.99 | 1.03 | 0.35 | 0.37 | M6 |
| AE Ara.............. | PC 18 | 174104.9 | -470320 | 6.48 | 1.01 | 0.30 | 0.23: | M2 |
| AS 255 | Hen 1525 | 175708.8 | -351534 | 8.25 | 0.88 | 0.23 | ... | K3 |
| AS 270 | Hen 1581 | 180533.7 | -20 2034 | 5.65 | 1.33 | 0.42 | 0.28 | M1 |
| AS 289 | Hen 1627 | 181222.2 | -1140 07 | 5.04 | 1.28 | 0.46 | 0.42 | M3.9 |
| Y CrA ${ }^{\text {a }} \ldots \ldots \ldots \ldots .$. | HD 166813 | 181422.9 | -425030 | 6.54 | 0.98 | 0.35 | ... | M5 |
| AS 304 | Hen 1691 | 182526.7 | -28 3557 | 7.40 | 0.88 | 0.31 | ... | M4 |
| AS 316 ${ }^{\text {a }} \ldots \ldots \ldots \ldots . . .$. | Hen 2-417 | 184232.8 | -211747 | 7.76 | ... | 0.31 | $\ldots$ | M |
| AS $327 . \ldots \ldots . . . . . .$. | Hen 1730 | 185316.7 | -24 2257 | 7.62 | 1.07 | 0.27 | $\ldots$ | M4 |
| CD -43 $14304 \ldots .$. | Hen 1924 | 210006.3 | -423850 | 7.60 | 0.90 | 0.18 | $\ldots$ | K5.5 |

[^1]narrow slit $\left(3^{\prime \prime}-5^{\prime \prime}\right)$, to obtain the best spectral resolution, and through a wide slit ( $\approx 10^{\prime \prime}$ ), to obtain accurate spectrophotometry. The standard stars were taken from Stone \& Baldwin (1983) for the Two Degree Field (2dF) observations, from Taylor (1984) for the 1991 CCD observations, and from Stone \& Baldwin (1983) as improved by Hamuy et al. (1992) for the 1995 CCD observations. The standard stars were always observed with the wide slit.

The data reductions were performed at the Cerro Tololo Inter-American Observatory (CTIO) La Serena Computing Center. All the observations were reduced to linear wavelength and flux scales using the program IRAF (Image Reduction Analysis Facilities). ${ }^{2}$ The spectra thus obtained were then measured using the task SPLOT in IRAF, on Sun

[^2]workstations available at the Centro de Procesamiento Digital de Imágenes (CPDI) of the Departamento de Astronomía de la Universidad de Chile; the blends were resolved using the DEBLEND option in the task SPLOT. When the spectra were obtained through both a narrow and a wide slit with the same equipment, the narrow spectra were reduced to the wide spectra without losing resolution; in this way, differences in the continuum or line intensities obtained in different epochs of observation cannot be attributed to differences in the slit widths.

The spectra of all the objects observed are presented in Figures 1-14.

The measured fluxes relative to $\mathrm{H} \beta$ are listed in Table 4; the last line of the table gives the logarithm of the observed flux in $\mathrm{H} \beta$. No individual mean internal errors are given for the different lines observed, but an analysis of the fluxes measured for the same object from different spectra obtained during the same night shows that they may be

TABLE 2
Instrumentation in the Different Observing Runs

| Run | Date | Telescope | Detector | Wavelength Range <br> (A) | FWHM <br> (A) | Filter |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1....... | 1987 Jul | 1.0 m | 2 dF | 3200-7100 | 5 | WG 360 |
| 2a...... | 1990 Jun | 1.5 m | GEC CCD | 3100-7600 | 17 | ... |
| 2b...... | 1990 Jun | 1.5 m | GEC CCD | 6100-10400 | 16 | RG 610 |
| 3a...... | 1991 Jun | 1.0 m | 2 dF | 3200-7100 | 5 | ... |
| 3b...... | 1991 Jun | 1.0 m | 2 dF | 4300-7800 | 5 | GG 455 |
| 4a...... | 1991 Jul | 1.5 m | GEC CCD | 3200-7700 | 17 | ... |
| 4b...... | 1991 Jul | 1.5 m | GEC CCD | 6100-10400 | 16 | RG 610 |
| 5a...... | 1995 Apr | 1.5 m | GEC CCD | 4000-5400 | 5 | ... ${ }^{\text {a }}$ |
| 5b...... | 1995 Apr | 1.5 m | GEC CCD | 4600-7100 | 9 | GG 455 |

TABLE 3
Observing Runs for the Different Objects

| ОbJect | RuN |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\frac{1987}{1}$ | 1990 |  | 1991 Jun |  | 1991 Jul |  | 1995 |  |
|  |  | 2a | 2b | 3 a | 3b | 4a | 4b | 5a | 5b |
| Hen $828 . \ldots . . . . . . . .$. | + |  |  |  |  | + | + |  |  |
| Hen $905 \ldots . . . . . . . .$. | + |  |  |  |  |  |  |  |  |
| Hen 1103............ | + |  |  |  |  | + | + |  |  |
| Hen 1213............ | + |  |  | + | + |  |  | + | + |
| KX $\operatorname{TrA} \ldots \ldots . . . . . .$. | + |  |  |  |  |  |  |  |  |
| AE Ara.............. |  |  |  |  |  | + | + |  |  |
| AS $255 . . . . . . . . . . .$. | $+$ |  |  |  |  | + | $+$ |  |  |
| AS $270 \ldots \ldots . . . . . .$. | + |  |  |  |  | + | + | + | + |
| AS 289 ............... | $+$ |  |  | $+$ | + |  |  |  |  |
| Y CrA ............... |  | + | + |  |  |  |  |  |  |
| AS $304 \ldots \ldots . . . . . .$. | + |  |  |  |  |  |  |  |  |
| AS $316 \ldots . . . . . . . . .$. | $+$ |  |  |  |  | $+$ |  |  |  |
| AS $327 \ldots \ldots . . . . . .$. | + |  |  |  |  | + |  |  |  |
| CD $-43^{\circ} 14304 \ldots .$. | $+$ |  |  | $+$ | + | $+$ | $+$ |  |  |

estimated as $\leq 10 \%$ for well-separated lines with $\log F(\lambda)>-12$, increasing up to about $50 \%$ for blended lines or for lines with $\log F(\lambda) \leq-14$. The error determined for $\log F(\mathrm{H} \beta)$ is in all cases of the order of 0.1 dex. The true errors may be quite larger, considering possible errors in the calibrations and in the atmospheric extinction curves adopted and the difficulties in the tracing of the continuum due to the noise and to the presence of absorption bands.

For the determination of the physical parameters deduced from these fluxes, the lines of the Balmer series have been corrected for the contamination produced by the


Fig. 1.-Spectra of Hen 828. Here and in all the figures that follow the spectra obtained at the earliest date are at the bottom, and those obtained at the latest date are at the top. The ordinates and the separation steps are in ergs $\mathrm{cm}^{-2} \mathrm{~s}^{-1}$; the spectra corresponding to 1991 are combined spectra, including the red and blue observations. Here the separation step is $4 \times 10^{-14}$.

Pickering series of He II. Besides, $\mathrm{H} \alpha$ has been corrected for the contamination produced by He ir Pf 14.

## 3. THE LATE COMPONENTS

The spectral types of the late components were obtained-assuming that they are all giants-by the $[\mathrm{TiO}]_{1},[\mathrm{TiO}]_{2}$, and [VO] indices given by Kenyon \& Fernández-Castro (1987) and by the [TiO] index given by Kenyon (1986); the indexes presented by Terndrup, Frogel, \& Whitford (1990) were also used. The determinations were made separately for each period of observations. When the spectral type seemed to be dependent of the wavelength of the features used, the latest spectral type obtained for each object was adopted, because of the possibility that the


Fig. 2.-Spectrum of Hen 905 obtained in 1987


Fig. 3.-Spectra of Hen 1103. Notice the intensity of the Balmer continuum in emission in 1991. The separation step is $4 \times 10^{-14}$.
absorption bands are being veiled in some cases by the nebular component. The final spectral type adopted was confirmed by visual comparison with the digital spectra published by Turnshek et al. (1985). In general, the differences between the types obtained for a given object in different periods are within what may be considered the errors of the observations, and the average of the different classifications has been adopted.

The knowledge of the spectral types allows an estimate of the distance of the different objects. The absolute $K$ magnitudes used for the $M$ giants were obtained from the calibration of the $M$ giants located in the Galactic bulge made by Frogel \& Whitford (1987), using a distance to the center of the Galaxy equal to 8.3 kpc . The observed $K$ magnitudes were taken from Table 1, and the corresponding $A_{K}$ absorptions were taken from Table 6 of the next section.


Fig. 4.-Spectra of Hen 1213. The separation step is $1 \times 10^{-13}$. In 1991 there is a hint that the Balmer continuum is in emission. All the observations are made with the same slit width, which allows a direct comparison of the lines intensities. In the 1995 spectrum, the greater strength of the emission lines of the Balmer series, is noticeable.


Fig. 5.-High-excitation spectrum of $\mathrm{KX} \operatorname{TrA}(\mathrm{Cn} 1-2)$ obtained in 1987.


Fig. 6.-Combined spectrum of AE Ara obtained in 1991. The Balmer continuum is in emission.


Fig. 7.-Spectra of AS 255. The separation step is $2 \times 10^{-13}$. The emission lines are more intense in 1987 than in 1991. In 1991 the Balmer continuum can be see faintly in emission.


Fig. 8.-Spectra of AS 270. The separation step is $5 \times 10^{-14}$. The difference in the intensity of the emission lines, mainly of the Balmer lines, is noticeable.


Fig. 9.-Spectra of AS 289. The separation step is $2 \times 10^{-13}$. As in Fig. 4 , both spectra were obtained with the same slit width. The emission lines are much more intense in 1987 than in 1991. Notice the strength and width of $\lambda 6825$.


FIG. 10.-Combined spectrum of Y CrA , obtained in 1990


Fig. 11.-Spectrum of AS 304 obtained in 1987


Fig. 12.-Spectra of AS 316. The separation step is $2 \times 10^{-13}$.


Fig. 13.-Spectra of AS 327. The separation step is $1 \times 10^{-13}$. In spite of the difference in the slit widths for the two spectra presented here, the difference in the strength of the He lines, mainly of $\lambda 5876$, is noticeable.


Fig. 14.-Spectra of CD $-43^{\circ} 14304$. Notice the Balmer discontinuity in emission.

The results are listed in Table 5, which gives, by column, (1) the name of the object, (2) the spectral type with its mean error, (3) the spectral types listed by other authors, (4) the distances obtained in this paper, and (5) the distances listed by other authors.

The internal errors of the distances corresponding to our internal errors in the spectral types and in the reddenings are small, running from 0.1 to 0.4 kpc . But the main errors originate in the adequacy of the calibration used and in the accuracy of the calibration, which we cannot evaluate.

## 4. THE NEBULAR COMPONENT

### 4.1. Reddening, Optical Depth, and Related Parameters

The reddening constants $\mathrm{C} \beta$ and the optical depths at $\mathrm{H} \alpha, \tau(\alpha)$, with their standard errors, have been determined simultaneously by the method presented by GutiérrezMoreno \& Moreno (1996), using the reddening curve published by Seaton (1979). To obtain the absolute $\mathrm{H} \beta$ intensity, the observed flux was corrected by reddening and by self-absorption. This last correction was determined assuming that all the $\mathrm{H} \beta$ radiation comes from the nebula, without contribution from the core, and applying the numerical procedure suggested as a good first approximation by Ahern (1975).

The results are listed in Table 6, which gives, by column, (1) and (2) the name and year of the observation of the object; (3) $\log F(\mathrm{H} \beta$ ), corrected by the contamination produced by the Pickering series of He II ; (4) the reddening constant at $\mathrm{H} \beta, \mathrm{C} \beta$, with its standard error; (5) the mean value of $\mathrm{C} \beta$ for each object, adopted in all the calculations that follow; (6) the optical depth at $\mathrm{H} \alpha$, also with its standard error; when a negative value-always very close to 0 -has been obtained for $\tau(\alpha)$, the value 0 has been adopted; (7) the reddening $E(B-V)=\mathrm{C} \beta / 1.47$; (8) the reddenings
$E(B-V)$ given by other authors; (9) the value of $A_{K}=$ $0.37 E(B-V)$ derived from (7); (10) the correction factor for self-absorption, $C_{s}$, with the standard errors derived from the observational errors; and (11) the logarithm of the absolute $\mathrm{H} \beta$ intensity, $\log I(\mathrm{H} \beta)$, which includes the corrections for interstellar reddening and for self-absorption; standard errors are also listed.

As could be expected, and was already mentioned in a previous paper (Gutiérrez-Moreno et al. 1997), $\tau(\alpha)$ and $C_{s}$ are well correlated. We have found a linear relation,

$$
C_{s}=1.023+0.078 \tau(\alpha), \quad \sigma_{C_{s}}= \pm 0.016
$$

that allows a rapid estimate of $C_{s}$ when $\tau(\alpha)$ is known.

### 4.2. Electron Temperatures and Densities

Because of the characteristics of the spectra of the objects we are studying, not many of the lines generally used for plasma diagnostics have been measured; the only line intensity ratio we have available is the nebular to auroral ratio of O iII, $I(5007+4959) / I(4363)$. Since this ratio alone does not permit the determination of $T_{e}$ and $N_{e}$, we have made an estimate of the electron temperatures and densities using the method given by Ferland \& Shields (1978), which involves the lines of [O III] $\lambda \lambda 4959,5007,4363$, [Ne III] $\lambda 3869$, and He I $\lambda 5876$. The use of their method allows the determination not only of the electron temperature and density of the nebula but also of the relative abundances $N\left(\mathrm{O}_{\text {III }}\right) / N\left(\mathrm{Ne}\right.$ III) and $N\left(\mathrm{O}_{\text {III }}\right) / N\left(\mathrm{He}_{\text {II }}\right)$. Nevertheless, we have observations of [ Ne III] $\lambda 3869$ in only a few cases, and we cannot combine the equations as proposed by Ferland \& Shields (1978); thus we have used, in addition to the equation for [O III], only the equation involving the relation for $I(5007+4959) / I(5876)$. The use of this last relation implies the knowledge of the relative abundance $A=N(\mathrm{O}$ III) $/ N(\mathrm{He}$ II), which is not known in advance. Thus, we have adopted as an initial value the abundance $N(\mathrm{O}) / N(\mathrm{He})$ found in the Sun (taken from de Freitas Pacheco et al. 1991): $A=8.685 \times 10^{-3}$. This assumes that $\mathrm{O}_{\mathrm{III}}$ and He II are present in nearly coincident regions and that their relative abundances are equal to that of the corresponding ions, which is adequate since their ionization potentials are similar; then

$$
\frac{N\left(\mathrm{O}_{\mathrm{III}}\right)}{N\left(\mathrm{He}_{\mathrm{II}}\right)}=\frac{N(\mathrm{O})}{N(\mathrm{He})}
$$

Besides, the method assumes Case $B$ conditions for the region where the He I lines are formed.

The numerical coefficients included in relations (1) and (3) of Ferland \& Shields (1978) have been modified according to the values of the atomic constants given by Mendoza (1983). We have found values of $T_{e}$ and $N_{e}$ with the adopted
TABLE 4
Observed Fluxes $(\mathrm{H} \beta+\operatorname{Pi} 8=100)$

| $\begin{gathered} \text { LiNe } \\ (\AA) \end{gathered}$ | Identification | Hen 828 |  | Hen 905 | Hen 1103 |  | Hen 1213 |  |  | $\frac{\text { KX TrA }}{1987}$ | AE Ara <br> 1991 | AS 255 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 1987 | 1991 | 1987 | 1987 | 1991 | 1987 | 1991 | 1995 |  |  | 1987 | 1991 |
| 3133 .......... | O III | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 3.73 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 3312 | O III | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 1.91 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |  |
| 3335 .......... | Ne II | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 6.53 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ |
| 3346 | [ Ne v] | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | ... | $\ldots$ | $\ldots$ | ... | 7.49 | $\ldots$ | ... |
| 3708 .......... | O III | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 8.90 | $\ldots$ | $\ldots$ |
| 3758 .......... | He II | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\cdots$ | $\ldots$ | 8.36 | $\ldots$ | $\ldots$ |
| 3796 | He II | ... | $\ldots$ | $\ldots$ | ... | 2.61 | ... | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | ... |
| 3835 ........... | He II +H 9 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 4.74 | 3.88 | 4.62 | $\ldots$ |  | 8.55 | 6.75 | $\ldots$ |
| 3869 | $\left[\mathrm{Ne}_{\mathrm{III}}\right]+\mathrm{He}_{\mathrm{I}}$ | 6.69 | $\ldots$ | ... |  | ... | 3.62 | 2.07 | $\ldots$ | 6.90 | 27.8 |  |  |
| 3889 ........... | $\mathrm{He}_{\mathrm{I}}+\mathrm{H} 8$ | 11.7 | ... | 10.6 | 9.55 | 11.9 | 6.59 | 6.70 | $\ldots$ | 10.3 | 14.4 | 5.14 | 11.5 |
| 3968 ........... | $[\mathrm{Ne} \mathrm{III]}]+\mathrm{H} 7$ | 13.4 | $\ldots$ | 13.4 | 10.4 | 7.87 | 8.56 | 8.25 | $\ldots$ | 10.5 | 26.3 | 10.2 | 8.87 |
| 4026 ........... | $\mathrm{He} \mathrm{II}^{\text {II }}$ | ... | $\ldots$ | ... | ... | ... | ... | ... | $\ldots$ | 1.74 | $\ldots$ | ... | ... |
| 4072 | [ $\mathrm{S}_{\text {II] }}$ ] | $\ldots$ | $\ldots$ | $\ldots$ |  | $\ldots$ | $\ldots$ |  |  | 1.20 | $\ldots$ | $\ldots$ |  |
| 4101 | N III $+\mathrm{H} \delta$ | 16.5 | 15.6 | 15.9 | 17.7 | 17.1 | 13.6 | 16.0 | 16.7 | 17.9 | 26.3 | 12.7 | 16.1 |
| 4200 | $\mathrm{He}_{\text {II }}$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | 1.51 | $\ldots$ | $\ldots$ | $\ldots$ |
| 4340 .......... | H $\gamma$ | 30.2 | 26.5 | 30.3 | 35.1 | 28.7 | 31.8 | 33.3 | 34.7 | 35.4 | 40.9 | 31.0 | 29.5 |
| 4363 | [ OIII ] $^{\text {l }}$ | 10.3 | 14.4 | ... | 3.51 | 4.83 | 5.57 | 4.64 | 3.98 | 8.13 | 30.8 | $\ldots$ | ... |
| 4388 | He I | ... | ... | $\ldots$ | ... | ... | 5.88 | ... | 4.12 | ... | ... | 3.49 | $\ldots$ |
| 4471 | He I | $\ldots$ | $\ldots$ | $\ldots$ | 2.98 | 3.35 | 3.38 | 3.90 | 2.83 | 2.13 | 7.00 | 2.14 | $\ldots$ |
| 4542 | He II | ... | $\ldots$ | $\ldots$ | ... | ... | ... | ... | ... | 2.86 | ... | ... | $\ldots$ |
| 4640 | N III | 9.56 | $\ldots$ | $\ldots$ | $\cdots$ | 18.3 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 4647 | C III |  |  |  |  | $\ldots$ | $\ldots$ |  | $\ldots$ |  | 16.1 |  |  |
| 4686 | He II | 30.9 | 13.8 | 98.9 | 82.2 | 31.5 | 18.4 | 16.6 | 9.12 | 59.7 | 20.6 | 95.4 | 97.0 |
| 4712 .......... | $\left[\mathrm{Ar} \mathrm{IV}^{\text {] }}\right.$ ] $\mathrm{He}_{\text {I }}$ | $\cdots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 8.14 | 5.28 | 3.83 | $\ldots$ | $\ldots$ | 3.04 |  |
| 4861 | $\mathrm{H} \beta$ | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 |
| 4922 | $\mathrm{He}_{\mathrm{I}}$ | 9.87 | 17.7 | ... | 4.10 | 11.2 | 18.2 | 13.9 | 11.0 | 3.91 | 7.69 | 8.3 | 9.76 |
| 4945 | [ Fe VII ] | $\ldots$ | ... | $\ldots$ | ... | ... | ... | ... | ... | 5.32 | $\ldots$ | ... | ... |
| 4959 | [ $\mathrm{O}_{\text {III] }}$ ] | 8.36 | $\ldots$ | $\cdots$ | $\cdots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | 6.81 | 29.1 | $\ldots$ | $\ldots$ |
| 5007 | [ $\mathrm{O}_{\mathrm{III} \text { ] }}$ | 25.0 |  | $\ldots$ | $\ldots$ | 22.6 | 3.66 | 4.03 | 2.92 | 23.4 | 86.5 | ... | $\ldots$ |
| 5016 | He I | 6.99 | $20.7{ }^{\text {a }}$ | $\ldots$ | $\ldots$ | ... | 12.1 | 4.00 | 4.92 | 2.89 | ... | 9.41 | $\ldots$ |
| 5048 ........... | He I | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | .. |  | $\ldots$ | 2.93 | $\ldots$ |
| 5159 | [ Fe VII ] | 2.66 | $\ldots$ | 2.13 | $\ldots$ | $\ldots$ | 5.15 | $\ldots$ | 4.42 | 5.00 | $\ldots$ | ... | $\ldots$ |
| $5278 \ldots \ldots .$. | [ Fe VII ] | 2.56 | $\ldots$ | ... | $\ldots$ | $\ldots$ | ... | $\ldots$ | ... | 3.14 | $\ldots$ | $\ldots$ | $\ldots$ |
| 5309 | [ Cav v] | ... | $\ldots$ | . | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 2.05 | $\ldots$ | $\ldots$ | $\ldots$ |
| 5412 | He II | $\ldots$ | $\ldots$ | 14.2 | 7.49 | 5.85 | $\ldots$ | $\ldots$ | $\ldots$ | 6.42 | $\ldots$ | 9.35 | $\ldots$ |
| 5615 | [Ca VII] | $\cdots$ | $\ldots$ | ... | ... | ... | $\ldots$ | $\cdots$ | $\cdots$ | 5.98 | $\cdots$ | ... | $\cdots$ |
| 5721 | [Fe viI] | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 10.4 | $\ldots$ | ... | $\ldots$ |
| 5807 ........... | He II | 4.68 | $\ldots$ | $\ldots$ |  | $\ldots$ | $\ldots$ |  | $\ldots$ |  | 20.1 | $\ldots$ |  |
| 5876 | He I | 35.1 | 37.0 | 39.5 | 22.3 | 21.9 | 64.3 | 45.9 | 48.2 | 12.6 | 40.8 | 18.8 | 39.1 |
| 6087 .......... | $[\mathrm{Ca} \mathrm{v}]+[\mathrm{Fe} \mathrm{viI]}$ | 8.38 | $\ldots$ | 3.60 | ... | ... | ... | ... | ... | 16.6 | ... | ... | ... |
| $6311 \ldots \ldots . .$. | $\mathrm{He}_{\text {II }}+[\mathrm{S} \mathrm{III}]$ | . | $\cdots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 2.05 | $\ldots$ | $\cdots$ | $\ldots$ |
| 6563 .......... | $\mathrm{H} \alpha$ | 984 | 825 | 1224 | 767 | 548 | 831 | 729 | 587 | 943 | 562 | 806 | 1009 |
| 6678 | $\mathrm{He}_{\text {I }}$ | 25.1 | 24.3 | 33.9 | 24.5 | 21.8 | 88.5 | 42.5 | 58.6 | 6.38 | 15.7 | 25.6 | 42.6 |
| 6825 ........... | O vi |  | $\ldots$ | 119 | 54.2 | $\cdots$ | . | $\ldots$ | ... | 49.4 | $\ldots$ | 41.9 | 119 |
| 7065 .......... | He I | 39.0 | $\ldots$ | 46.0 | ... | 17.5 | 77.8 | 30.2 | $\ldots$ | 13.0 | 26.4 | 15.4 | 57.4 |
| 7281 .......... | He I | ... | $\ldots$ | ... | $\ldots$ | 5.57 | ... | ... | $\ldots$ | ... | ... | ... |  |
| 10124 .......... | He II | $\ldots$ | ... | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 40.8 |
| $-\log F(\mathrm{H} \beta) \ldots$ |  | 12.2 | 12.6 | 12.8 | 12.2 | 11.8 | 12.0 | 12.1 | 11.9 | 11.2 | 11.6 | 11.8 | 12.6 |

TABLE 4-Continued

| $\begin{gathered} \text { LINE } \\ (\AA) \end{gathered}$ | IdEntification | AS 270 |  |  | AS 289 |  | $\frac{\text { Y CrA }}{1990}$ | $\frac{\text { AS } 304}{1987}$ | AS 316 |  | AS 327 |  | CD $-43^{\circ} 14304$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 1987 | 1991 | 1995 | 1987 | 1991 |  |  | 1987 | 1991 | 1987 | 1991 | 1987 | 1991 |
| 3133 | O III | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 6.33 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 3203 .......... | He II | $\ldots$ | $\ldots$ | ... | ... | ... | ... | ... | $\ldots$ | 14.7 | $\ldots$ | 15.8 | ... | ... |
| 3341 ......... | $\mathrm{O}_{\text {III }}+[\mathrm{Ne} v]$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 89.8 | $\ldots$ | $\ldots$ | ... | $\ldots$ |  | $\ldots$ |  |
| 3423 | $\mathrm{O}_{\text {III }}+[\mathrm{Ne} \mathrm{v}]$ | $\ldots$ | $\ldots$ | ... | ... | $\ldots$ | 283 | $\ldots$ | $\ldots$ | $\ldots$ | ... | 22.5 | ... | 7.26 |
| 3429 | O III | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | 13.0 | $\ldots$ | ... | $\ldots$ | ... |
| 3499 | He I | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 7.07 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 3586 ......... | $[\mathrm{Fe} \mathrm{viI]}]+\mathrm{He}$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | 19.7 | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | ... |
| 3727 | [ $\mathrm{OII}^{\text {] }}$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | 11.3 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 3759 .......... | [ Fe VII ] | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 33.7 | $\ldots$ | $\ldots$ | ... |  | $\ldots$ | $\ldots$ | $\ldots$ |
| 3835 ......... | He II +H 9 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | 12.6 | 3.90 | 4.01 | 2.66 | $\ldots$ | $\ldots$ | $\ldots$ |
| 3869 .......... | $[\mathrm{Ne} \mathrm{III}]+\mathrm{He}_{\mathrm{I}}$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 32.8 | 8.92 | ... | ... | ... | ... | $\ldots$ | $\ldots$ |
| 3889 | $\mathrm{He}_{\mathrm{I}}+\mathrm{H} 8$ | ... | $\ldots$ | $\ldots$ | 3.62 | $\ldots$ | 15.1 | 22.7 | 7.40 | 8.00 | 5.40 | 9.60 | 5.96 | 11.3 |
| 3968 ........... | $[\mathrm{Ne}$ III] $]+\mathrm{H} 7$ | $\ldots$ | $\ldots$ | $\ldots$ | 5.47 | $\ldots$ | 25.8 | 24.5 | 6.17 | 7.87 | 5.98 | 9.93 | 4.84 | 8.80 |
| 4026 ........... | He II | ... | $\ldots$ |  | 2.87: | ... | ... | 5.62 | ... | ... | 3.53 | ... | 5.69 | 4.65 |
| 4101 | N III $+\mathrm{H} \delta$ | 13.3 | 15.3 | 13.5 | 10.6 | 10.4 | 24.6 | 23.3 | 14.6 | 14.7 | 13.9 | 15.8 | 13.9 | 15.3 |
| 4200 .......... | He II |  |  |  | ... |  | ... | $\ldots$ | ... |  | 3.32 | ... |  | ... |
| 4340 | $\mathrm{H} \gamma$ | 27.8 | 30.5 | 24.3 | 24.2 | 23.5 | 44.3 | 44.4 | 32.7 | 31.5 | 29.3 | 33.7 | 29.7 | 33.0 |
| 4363 .......... | [ OIII ] | 16.5 | 15.3 | 6.68 | 1.38 | 4.34 | 32.0 | 5.68 | ... | ... | ... | ... | 10.7 | 4.70 |
| 4388 | He I | 5.91 | ... | 2.00 | ... | ... | ... | 6.51 | $\ldots$ | $\ldots$ | 3.96 | ... | ... | 5.35 |
| 4415 | [ $\mathrm{O}_{\mathrm{II} \text { ] }}$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | 3.94 | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... |
| 4471 | $\mathrm{He}_{\text {I }}$ | 6.76 | 6.92 | 1.97 | 2.67 | $\ldots$ | $\ldots$ | 5.21 | 2.02 | 2.34 | 3.49 | 3.00 | 4.16 | 8.73 |
| 4542 | He II | ... | ... | ... | ... | $\ldots$ | $\ldots$ | ... | 2.17 | 2.16 | 5.52 | ... | ... | ... |
| 4647 .......... | C III | $\ldots$ | 6.92 | $\ldots$ |  |  | $\ldots$ | $\ldots$ | $\ldots$ |  | . | .. |  |  |
| 4686 | He II | 41.6 | 45.3 | 23.1 | 64.1 | 82.3 | 94.0 | 23.5 | 52.6 | 47.9 | 91.5 | 73.2 | 65.2 | 68.9 |
| 4712 .......... | [ Ar IV ] $+\mathrm{He}_{\mathrm{I}}$ | 8.73 | 20.5 | 4.29 | ... | ... | ... | 4.02 | ... | ... | 3.41 | ... | ... | 6.54 |
| 4716 | [ Ne IV] |  | $\ldots$ |  | $\ldots$ | $\ldots$ | 9.12 | ... | $\ldots$ | $\ldots$ | ... | $\cdots$ | $\cdots$ | ... |
| 4725 ........... | [ $\mathrm{Ne}_{\text {IV }}$ ] | 11.7 | $\ldots$ | 3.56 | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\cdots$ | $\ldots$ | $\ldots$ |
| 4740 | [ Ar IV ] |  |  |  |  |  | 11.5 |  |  |  |  |  |  |  |
| 4861 .......... | $\mathrm{H} \beta$ | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 | 100 |
| 4922 .......... | He I | 18.6 | 64.6 | 5.89 | 8.38 | ... | $\ldots$ | 15.0 | 6.18 | 10.1 | 6.14 | 11.1 | 3.79 | 12.0 |
| 5007 | [O III] | 30.4 | 26.2 | 9.84 | 2.18 | ... | 103 | 9.16 | 1.17: | ... | 3.02: | ... |  |  |
| 5016 ........... | He I | 9.63 | ... | 5.11 | 8.69 | $\ldots$ | ... | 14.4 | 4.63 | 4.64 | 8.68 | 2.62 | 8.49 | 10.1 |
| 5159 ........... | [Fe viI] | ... | $\ldots$ | $\ldots$ | 4.85: | $\ldots$ |  | ... | ... | $\ldots$ | 4.78 | ... | ... | ... |
| 5309 | [Ca v] | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 10.5 | $\ldots$ | ... | $\ldots$ | ... | $\cdots$ | $\ldots$ | $\ldots$ |
| 5412 | He II | $\ldots$ | $\ldots$ | $\ldots$ | 12.9 | 28.7 | 27.2 | $\ldots$ | 7.16 | 6.29 | 10.9 | 10.1 | $\ldots$ | 8.47 |
| 5721 | [Fe viI] | $\ldots$ | $\ldots$ |  | 11.2 | 20.5 | 31.3 | $\ldots$ | $\ldots$ |  | 4.05 | $\ldots$ |  | ... |
| 5876 | He I | 111 | 125 | 51.5 | 54.9 | 62.5 | 28.8 | 31.7 | 22.0 | 15.5 | 50.1 | 17.2 | 50.1 | 52.1 |
| 6087 | $[\mathrm{Ca} \mathrm{v}]+[\mathrm{Fe} \mathrm{viI]}$ | $\ldots$ | ... | ... | 26.7 | 45.5 | 45.6 | ... | ... | ... | 6.10 | ... | ... | ... |
| 6528 .......... | He II | ... | ... | ... | ... | ... | 45.9 | $\ldots$ | ... | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ |
| 6563 | $\mathrm{H} \alpha$ | 1922 | 2325 | 1514 | 2152 | 2228 | 679 | 659 | 938 | 717 | 1133 | 903 | 897 | 631 |
| 6678 ........... | He I | 115 | 160 | 79.2 | 91.6 | 114 | 12.8 | 31.8 | 18.3 | 14.4 | 58.8 | 23.0 | 58.0 | 28.7 |
| 6825 | O VI | $\ldots$ | ... | ... | 67.9 | 331 | 53.2 | ... | 87.9 | 83.7 | 119 | 92.6 | 141 | 75.3 |
| 7065 .......... | He I | 164 | ... | ... | 83.6 | 75.2 | ... | 24.4 | 20.5 | 8.57 | 63.4 | 15.8: | 37.6 | 38.6 |
| 9232 .......... | $\mathrm{HI}_{\mathrm{I}}(\mathrm{Pa} 9)$ | ... | $\ldots$ | $\ldots$ | ... | ... | $\ldots$ | ... | ... | ... | ... | ... | ... | 15.6 |
| $9403 \ldots \ldots . .$. | ? | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 32.4 |
| 10049 .......... | $\mathrm{HI}_{( }(\mathrm{Pa} \delta)$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 10.8 |
| 10124 ......... | He II |  |  |  |  |  |  |  |  |  | $\ldots$ | $\ldots$ | $\ldots$ | 22.6 |
| $-\log F(\mathrm{H} \beta) \ldots$ |  | 13.0 | 13.0 | 12.5 | 12.0 | 12.7 | 12.9 | 11.5 | 11.8 | 11.8 | 12.3 | 12.2 | 12.1 | 11.6 |

[^3]TABLE 5
Spectral Type of the Late Component and Derived Parameters

| Object <br> (1) | Spectral Type |  | $\begin{gathered} \text { Distance } \\ (\mathrm{kpc}) \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | This Paper <br> (2) | Other <br> (3) | This Paper <br> (4) | Other <br> (5) |
| Hen $828 . . . . . . . . . . . .$. | $\mathrm{M} 2.9 \pm 0.4$ | M5.5 ${ }^{\text {a }}$ | 1.6 | $2.2^{\text {a }}$ |
| Hen $905 \ldots . . . . . . . . .$. | M4.3 $\pm 0.1$ | M $2^{\text {a }}$ | 3.9 | $1.7{ }^{\text {a }}$ |
| Hen 1103............ | $\mathrm{M} 1.3 \pm 0.7$ |  | 2.3 |  |
| Hen 1213............ | $\mathrm{M} 0.0 \pm 0.2$ |  | 0.9 |  |
| KX TrA ............. | $\mathrm{M} 3.9 \pm 0.6$ | M5.5, ${ }^{\text {a }}$ M3 ${ }^{\text {b }}$ | 1.2 | 3.8, ${ }^{\text {a }} 3.0{ }^{\text {b }} 4.4{ }^{\text {c }}$ |
| AE Ara.............. | M4.4 $\pm 0.3$ | M3 ${ }^{\text {b }}$ | 1.9 | $2.3,{ }^{\text {b }} 1.6{ }^{\text {c }}$ |
| AS $255 \ldots \ldots . . . . .$. | $\mathrm{K} 5.2 \pm 0.2$ |  | 1.5 |  |
| AS $270 \ldots . . . . . . . . .$. | M4.3 $\pm 0.2$ |  | 1.1 |  |
| AS 289 .............. | M3.3 $\pm 0.4$ |  | 0.6 |  |
| Y CrA ............... | M6.1 $\pm 0.1$ |  | 2.9 | $3.9{ }^{\text {c }}$ |
| AS 304 | $\mathrm{M} 0.1 \pm 0.8$ |  | 1.3 | $2.5{ }^{\text {d }}$ |
| AS 316 | $\mathrm{M} 2.7 \pm 0.2$ | M3 ${ }^{\text {b }}$ | 1.7 | $4.5{ }^{\text {b }}$ |
| AS 327 | $\mathrm{M} 2.5 \pm 0.4$ | $\mathrm{M} 0^{\text {b }}$ | 1.9 | $3.0{ }^{\text {b }}$ |
| CD -43 $14304 \ldots \ldots$ | $\mathrm{M} 0.3 \pm 0.1$ | M0, ${ }^{\text {a }}$ K $5.5{ }^{\text {e }}$ | 1.3 | $1.0{ }^{\text {a }}$ |

[^4]abundance ratio for all the objects analyzed, except for the 1987 observations of Hen 1213, for which we got solutions only for $A \leq 4.9 \times 10^{-3}$. A second approximation made using the ratio $A$ determined with the values of $T_{e}$ and $N_{e}$ just obtained showed only small variations.

We have to consider that the values of $T_{e}$ and $N_{e}$ obtained by this method, although better than values arbitrarily adopted, must be considered only as rough estimates. In fact, in addition to the observational errors in the data, which in some cases may be large, and the possible error involved in the a priori adoption of a relative abundance $N(\mathrm{O}$ III) $/ N(\mathrm{He}$ II), which may not be adequate, we have two sources of error that may be important: (a) in S-type symbiotic nebulae the He I lines may not be formed in the same region than the [ O miI ] lines, and (b) Proga, Mikolajewska, \& Kenyon (1994) have found that the assumption of Case B conditions for the region where $\mathrm{He}_{\mathrm{I}}$ lines are formed is not true for all S-type symbiotic systems.

For objects where [O III] $\lambda 5007$ is too faint to be measured, we have adopted approximate temperatures $T_{e}=$ $15,000 \mathrm{~K}$; in these cases, the values of $N_{e}$ have been estimated from the diagnostic diagram by Gutierrez-Moreno et al. (1995) as being of the order of $10^{9.5} \mathrm{~cm}^{-3}$. If an object was observed in two or more different runs and observations of [O III] $\lambda 5007$ were available in only one of them, the values of $T_{e}$ and $N_{e}$ obtained in that run were adopted for the other ones.

Our values of $T_{e}$ and $N_{e}$, together with values published by other authors, are listed in columns (3)-(6) of Table 7.

### 4.3. Helium Abundance

Owing to the high density, large self-absorption corrections are expected. Collisional excitation corrections were performed following Peimbert \& Torres-Peimbert (1987) before computing the He II abundances by the procedure followed by Gutiérrez-Moreno \& Moreno (1996) and Gutiérrez-Moreno et al. (1997). The lines used were He I $\lambda \lambda 4471,5876$, and 7065; the emissivities for these lines were interpolated for the different values of $T_{e}$ and $N_{e}$ from the data computed by Almog \& Netzer (1989) for $N_{e}=10^{4}$, $10^{6}, 10^{8}$, and $10^{10}$ and $T_{e}=10,000 \mathrm{~K}$. He I $\lambda 6678$ was not used since few emissivities for this line are listed by Almog \& Netzer (1989) for these high densities. $\mathrm{H} \gamma$ was used as the reference line in order to diminish the effects of optical depth in the hydrogen lines, since we know that $\tau(\gamma) \approx 0.07 \tau(\alpha)$. When He I $\lambda 7065$ was not observed, we adopted for $\tau(3889)$ the values derived from previous observations of the same object, which may not be very adequate. For one object ( Y CrA ) we had only one observation, with no values for $\lambda 7065$; thus, to have an estimate of He abundance we adopted the value of $\tau(3889)$ obtained for AE Ara, which has similar values of $T_{e}$ and $N_{e}$. The resulting error in

TABLE 6
Reddening and Related Parameters

${ }^{\text {a }}$ Feibelman 1991 (UV values).
${ }^{\text {b }}$ Pereira 1995.
${ }^{\text {c }}$ Mikolajewska et al. 1997.
${ }^{\text {d }}$ Nussbaumer et al. 1988 (UV values).
${ }^{\text {e }}$ Munari \& Buson 1993.
${ }^{\mathrm{f}}$ Schmid \& Nussbaumer 1993 (UV values).
$y^{+}$must not be large, since changes in the value of $\tau(3889)$ do not produce large variations in the emissivity of $\lambda 5876$.

The He III abundance was computed from He II $\lambda 4686$, using Brocklehurst (1971) emissivities extrapolated to our values of the density.

The results on He abundances are listed in columns (7)(12) of Table 7, which give (7) the $\mathrm{He}^{+}$abundance, expressed as $y^{+}=N\left(\mathrm{He}^{+}\right) / N\left(\mathrm{H}^{+}\right)$, obtained as the direct average of the lines used except when $I(7065)$ was not measured, in which case double weight was given to $y^{+}(5876)$; (8) the value of $y^{++}=N\left(\mathrm{He}^{++}\right) / N\left(\mathrm{H}^{+}\right)$; (9) the total He abundance, $y=y^{+}+y^{++}=N(\mathrm{He}) / N(\mathrm{H})$. Columns (10), (11), and (12) list the same values as listed by Pereira (1995). The errors in $y^{+}$and $y^{++}$are given in units of the last decimal; since we have used only one line for the determination of $y^{++}$, the error quoted in this case is an estimate based on the mean error of the measures of this line. Since the values of $T_{e}$ and $N_{e}$ are only estimates, the real errors may be quite larger than those quoted here.

Comparison of our values with Pereira's (1995) for the objects in common shows large differences in $y^{+}, y^{++}$, and
$y$; they may be due to the differences in reddening and the consequent differences in $T_{e}$ and $N_{e}$, but they are large enough to suspect that some variability may be present. In fact, treating Pereira's (1995) data with our procedures gives slightly larger He abundances, but the differences with our values are still too large to be explained only by the observational errors.

## 5. THE HOT SOURCES

### 5.1. Magnitudes

$B$ and $V$ magnitudes were obtained for all objects observed by convolving the intensities measured in each spectrum-without considering the emission lines due to the nebula-with the sensitivity functions of Johnson's $B$ and $V$ filters. The results are listed in Table 8, which gives the magnitudes thus determined and the corresponding $B-V$ colors; the errors listed for $B$ and $V$ are the mean errors of the average of the magnitudes obtained for each object and epoch. The table includes also the values listed

TABLE 7
Nebular Parameters

| Object <br> (1) | Year <br> (2) | $T_{e}$ |  | $\log N_{e}$ |  | This Paper |  |  | Other |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | This <br> Paper <br> (3) | Other <br> (4) | This <br> Paper <br> (5) | Other <br> (6) | $y^{+}$ <br> (7) | $y^{++}$ <br> (8) | $\begin{gathered} y \\ (9) \end{gathered}$ | $\begin{gathered} y^{+} \\ (10) \end{gathered}$ | $y^{++}$ <br> (11) | $\begin{gathered} y \\ (12) \end{gathered}$ |
| Hen 828 | 1987 | 10700 | $\ldots$ | 7.9 | $\ldots$ | $0.079 \pm 0.001$ | $0.037 \pm 0.001$ | 0.116 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 10700 | $\ldots$ | 7.9 | $\ldots$ | $0.080 \pm 0.011$ | $0.019 \pm 0.003$ | 0.099 | $\ldots$ | $\ldots$ | $\ldots$ |
| Hen $905 \ldots . . . . . . . . .$. | 1987 | 15000 | $\ldots$ | 9.5 | $\ldots$ | $0.061 \pm 0.002$ | $0.142 \pm 0.001$ | 0.203 | $\ldots$ | $\ldots$ | $\ldots$ |
| Hen 1103............ | 1987 | 9300 | $\ldots$ | 7.6 | $\ldots$ | $0.063 \pm 0.001$ | $0.088 \pm 0.004$ | 0.151 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 9300 | ... | 7.6 | $\ldots$ | $0.085 \pm 0.004$ | $0.040 \pm 0.001$ | 0.125 | $\ldots$ | $\ldots$ | $\ldots$ |
| Hen 1213............ | 1987 | 15700 | $\ldots$ | 9.1 | $\ldots$ | $0.099 \pm 0.006$ | $0.024 \pm 0.002$ | 0.123 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 13800 | ... | 9.1 | $\ldots$ | $0.091 \pm 0.003$ | $0.014 \pm 0.001$ | 0.105 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1995 | 14900 | , | 9.3 | , | $0.082 \pm 0.010$ | $0.011 \pm 0.003$ | 0.093 | $\ldots$ | ... | , |
| KX $\operatorname{TrA} \ldots \ldots \ldots \ldots .$. | 1987 | 10500 | $16000^{\text {a }}$ | 7.7 | $6.4{ }^{\text {a }}$ | $0.037 \pm 0.005$ | $0.062 \pm 0.003$ | 0.099 | $0.64{ }^{\text {a }}$ | $0.036{ }^{\text {a }}$ | $0.100^{\text {a }}$ |
| AE Ara.............. | 1991 | 10000 | $10500^{\text {a }}$ | 7.5 | $7.0^{\text {a }}$ | $0.161 \pm 0.005$ | $0.021 \pm 0.001$ | 0.182 | $0.80{ }^{\text {a }}$ | $0.012^{\text {a }}$ | $0.092^{\text {a }}$ |
| AS 255 | 1987 | 15000 | ... | 9.5 | ... | $0.036 \pm 0.004$ | $0.136 \pm 0.008$ | 0.172 | ... | ... | ... |
|  | 1991 | 15000 | $\ldots$ | 9.5 | $\ldots$ | $0.058 \pm 0.000$ | $0.145 \pm 0.001$ | 0.203 | $\ldots$ | $\ldots$ | $\ldots$ |
| AS 270 | 1987 | 11900 | $\ldots$ | 8.3 | $\ldots$ | $0.170 \pm 0.004$ | $0.051 \pm 0.005$ | 0.221 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 12100 | $\ldots$ | 8.4 | $\ldots$ | $0.163 \pm 0.011$ | $0.050 \pm 0.001$ | 0.213 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1995 | 12800 | $\ldots$ | 8.5 | $\ldots$ | $0.068 \pm 0.009$ | $0.033 \pm 0.003$ | 0.101 | $\ldots$ | $\ldots$ | $\ldots$ |
| AS 289 | 1987 | 13000 | $\ldots$ | 9.1 | ... | $0.058 \pm 0.002$ | $0.091 \pm 0.003$ | 0.149 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 13000 | ... | 9.1 | $\ldots$ | $0.073 \pm 0.000$ | $0.122 \pm 0.011$ | 0.195 | $\ldots$ | $\ldots$ | $\ldots$ |
| Y CrA ............... | 1990 | 10200 | $13000^{\text {a }}$ | 7.4 | $6.5^{\text {a }}$ | $0.103 \pm 0.005$ | $0.088 \pm 0.004$ | 0.191 | $0.048^{\text {a }}$ | $0.037^{\text {a }}$ | $0.085^{\text {a }}$ |
| AS 304 | 1987 | 10400 | ... | 8.4 | ... | $0.098 \pm 0.002$ | $0.024 \pm 0.003$ | 0.122 | ... | ... | ... |
| AS 316 | 1987 | 13400 | $\ldots$ | 9.2 | $\ldots$ | $0.038 \pm 0.000$ | $0.068 \pm 0.003$ | 0.106 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 13400 | $\ldots$ | 9.2 | $\ldots$ | $0.031 \pm 0.007$ | $0.065 \pm 0.002$ | 0.096 | $\ldots$ | $\ldots$ | $\ldots$ |
| AS 327 | 1987 | 17100 | $\ldots$ | 9.3 | $\ldots$ | $0.078 \pm 0.004$ | $0.129 \pm 0.001$ | 0.207 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | 17100 | $\ldots$ | 9.3 |  | $0.037 \pm 0.009$ | $0.088 \pm 0.002$ | 0.125 | $\cdots$ | $\ldots$ | $\ldots$ |
| CD $-43^{\circ} 14304 \ldots \ldots$ | 1987 | 15000 | $\ldots$ | 9.5 | $9{ }^{\text {b }}$ | $0.083 \pm 0.004$ | $0.094 \pm 0.004$ | 0.177 | $\ldots$ | ... | $\ldots$ |
|  | 1991 | 15000 | $\ldots$ | 9.5 | $\cdots$ | $0.108 \pm 0.001$ | $0.098 \pm 0.002$ | 0.206 | $\cdots$ | $\cdots$ | $\cdots$ |

${ }^{\text {a }}$ Pereira 1995.
${ }^{\text {b }}$ Schmid \& Nussbaumer 1993.
by Munari et al. (1992) for some of our objects. As could be expected by not including the emission lines, our magnitudes are in general fainter than theirs.

### 5.2. Temperature, Luminosities, and Radii

We have estimated the radiation temperature of the hot sources by the method proposed by Iijima (1981); the values obtained may be considered upper limits. The results are listed in Table 9, together with some values presented by other authors. We have estimated errors of $7 \%$ in the mean for our temperatures. Mikolajewska, Acker, \& Stenholm (1997) obtain the temperatures by the procedure suggested by Mürset \& Nussbaumer (1994), who use a simple formula that relates the temperature of the hot component with the highest ionization potential observed in the nebular spectrum.

The luminosity was estimated by the method of Zijlstra \& Pottasch (1989), using the line He II $\lambda 4686$ instead of $\mathrm{H} \beta$ :

$$
\frac{4 \pi d^{2} I\left(\mathrm{He}_{\mathrm{II}}\right)}{L^{*}}=h v(\mathrm{He} \mathrm{II}) \frac{\alpha\left(\mathrm{He}_{\mathrm{II}}\right)}{\alpha_{B}} \times \frac{15 G_{4}\left(T^{*}\right)}{\pi^{4} k T^{*}},
$$

where $I\left(\mathrm{He} \mathrm{II}^{\mathrm{II}}\right)$ is the He II $\lambda 4686$ flux corrected for reddening; $\alpha\left(\mathrm{He}_{\text {II }}\right)$ is the total recombination coefficient for He II $\lambda 4686$, whereas $\alpha_{B}$ is the effective recombination coefficient for the same line; and $G_{4}\left(T^{*}\right)$ is the Zanstra temperature integral for $\mathrm{He}_{\text {II }} \lambda 4686$. The values of $G_{4}\left(T^{*}\right), \alpha(\mathrm{He}$ II), and $\alpha_{B}$ have been tabulated (see, for example, Pottasch 1984). Knowing the luminosity, the radius may be easily estimated from the relation

$$
L^{*} / L_{\odot}=\left(R^{*} / R_{\odot}\right)^{2}\left(T^{*} / T_{\odot}\right)^{4} .
$$

The results obtained are listed in Table 9, which gives, for each object and each epoch of observation, in addition to the temperatures, the corresponding radii and luminosities and results obtained by other authors.

We have estimated errors of the order of $40 \%$ in the mean for the luminosities and radii, but the real errors may be much larger than these, since they not only depend on the errors of the temperatures quoted before, but also depend very strongly on the errors of the distances.

Since we have used the upper limit for $T^{*}$, the radius and luminosity listed are both lower limits.

TABLE 8
$B$ and $V$ Magnitudes

| Оbject | Year | This Paper |  |  | OTHER |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $B$ | $V$ | $B-V$ | $B$ | $V$ | $B-V$ |
| Hen 828 | 1987 | $15.52 \pm 0.02$ | $14.90 \pm 0.11$ | 0.62 | 15.19 | 14.29 | 0.90 |
|  | 1991 | $14.28 \pm 0.10$ | $14.08 \pm 0.06$ | 0.20 | ... | ... |  |
| Hen 905 | 1987 | $16.70 \pm 0.02$ | $15.43 \pm 0.04$ | 1.27 | 15.80 | 14.16 | 1.64 |
| Hen 1103............ | 1987 | $15.56 \pm 0.12$ | $14.52 \pm 0.03$ | 1.04 | 14.00 | 13.20 | 0.80 |
|  | 1991 | $13.05 \pm 0.01$ | $12.82 \pm 0.01$ | 0.23 | ... | $\ldots$ | $\ldots$ |
| Hen 1213............ | 1987 | $14.15 \pm 0.12$ | $12.75 \pm 0.06$ | 1.40 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1991 | $14.29 \pm 0.08$ | $12.81 \pm 0.08$ | 1.48 | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 1995 | $14.11 \pm 0.10$ | $12.79 \pm 0.04$ | 1.32 | $\ldots$ | $\ldots$ | $\ldots$ |
| KX $\operatorname{TrA} \ldots \ldots . . . . . .$. | 1987 | $14.78 \pm 0.11$ | $14.35 \pm 0.19$ | 0.43 | ... | $\ldots$ |  |
| AE Ara .............. | 1987 | $13.60 \pm 0.12$ | $13.48 \pm 0.06$ | 0.12 | 12.66 | 12.14 | 0.52 |
| AS 255 | 1987 | $14.47 \pm 0.04$ | $13.31 \pm 0.05$ | 1.16 | 14.75 | 13.01 | 1.74 |
|  | 1991 | $15.11 \pm 0.06$ | $13.67 \pm 0.01$ | 1.44 | ... | ... | ... |
| AS 270 | 1987 | $16.62 \pm 0.11$ | $14.75 \pm 0.11$ | 1.87 | 15.70 | 13.71 | 1.99 |
|  | 1991 | $16.29 \pm 0.08$ | $14.67 \pm 0.08$ | 1.62 | ... | ... | ... |
|  | 1995 | $15.96 \pm 0.08$ | $14.56 \pm 0.04$ | 1.40 | $\ldots$ | ... | $\ldots$ |
| AS 289 | 1987 | $15.72 \pm 0.19$ | $13.81 \pm 0.04$ | 1.91 | 15.72 | 13.62 | 2.10 |
|  | 1991 | $16.34 \pm 0.10$ | $14.35 \pm 0.04$ | 1.99 | ... | ... | ... |
| Y CrA ............... | 1990 | $15.50 \pm 0.12$ | $14.71 \pm 0.06$ | 0.79 | 15.56 | 14.35 | 1.21 |
| AS 304 | 1987 | $12.23 \pm 0.08$ | $12.06 \pm 0.08$ | 0.17 | 11.69 | 11.19 | 0.50 |
| AS 316 | 1987 | $15.12 \pm 0.05$ | $13.89 \pm 0.04$ | 1.23 | ... | ... | $\ldots$ |
|  | 1991 | $15.48 \pm 0.12$ | $14.41 \pm 0.08$ | 1.07 | ... | $\ldots$ |  |
| AS 327 | 1987 | $15.23 \pm 0.10$ | $13.99 \pm 0.02$ | 1.24 | 14.43 | 12.89 | 1.54 |
|  | 1991 | $15.30 \pm 0.12$ | $13.88 \pm 0.04$ | 1.42 | ... | ... |  |
| CD $-43^{\circ} 14304 \ldots \ldots$ | 1987 | $13.70 \pm 0.01$ | $12.53 \pm 0.04$ | 1.17 | 13.12 | 11.55 | 1.57 |
|  | 1991 | $13.19 \pm 0.08$ | $12.14 \pm 0.02$ | 1.05 | $\ldots$ | $\ldots$ | $\ldots$ |

Figure 15 presents the location of the objects observed in the $\mathrm{H}-\mathrm{R}$ diagram. For each object with several observations we have used the average of the temperatures and luminosities obtained, except in the cases of Hen 828 and Hen 1103 , for which the variations seem to be real. The objects indicated by a dot have He abundances greater than 0.17 .

## 6. DISCUSSION OF SOME INDIVIDUAL OBJECTS

### 6.1. Hen 828

There is a notorious increase of the level of the continuum between 1987 and 1991, and, as can be seen in Figure 1 and Table 4, many of the emission lines visible in 1987 have become undetectable in 1991, as if being veiled by the high level of the continuum; the lines that remain visible are mainly recombination lines; the only exceptions are the [ O III] lines $\lambda 4363$ and $\lambda 5007$, which are visible but have not been measured since they are completely blended with He I 25016. The dereddened 1991 spectrum (Fig. 16) shows a continuum with a maximum ( $\approx 3200-3400 \AA$ ), which suggests a temperature of the order of $7500-8500 \mathrm{~K}$, corresponding an A5-F0 star. These changes suggest that the object was having an outburst in 1991. Table 8 shows that

Hen 828 was brighter in both $B$ and $V$ during this outburst than in 1987; nevertheless, the resulting total luminosity was smaller for the 1991 observations, because of the faintness of He iI $\lambda 4686$.

### 6.2. Hen 905

In the second-order spectrum of Hen 905, an emission located at a distance of about $2^{\prime}$ to the west of the central object was detected in both $\mathrm{H} \alpha$ and O vi 26825 . The region between the object and this emission is full of faint emissions, which extend also to the east of the object, as shown in Figure 17, which corresponds to O vi 26825 . The minimum $X$-value of the plot is fixed at 0 and does not allow us to see the point symmetrical to the west emission; this fact suggests that there could be another emission at the east of the object, which would imply the existence of a ring surrounding the central object, and emitting mainly in $\mathrm{H} \alpha$ and O v. The emission in $\mathrm{H} \alpha(\approx 20)$ is more intense than in O vi $\lambda 6825(\approx 15)$ but is less conspicuous in the corresponding graph, since the intensity relative to the central line ( $\approx \frac{1}{6}$ ) is much smaller than in the case of $\lambda 6825$.
The emission in $\lambda 6825$ in the one-dimensional spectrum is fairly intense and wide, suggesting several components.

TABLE 9
Temperature, Radius, and Luminosity of the Нot Components

| Object | Year | $T^{*}$ |  | $\log L / L_{\odot}$ |  | $R / R_{\odot}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | This Paper | Other | This Paper | Other | This Paper | Other |
| Hen 828 | 1987 | 130,000 | $\ldots$ | 2.38 | $\ldots$ | 0.031 | $\ldots$ |
|  | 1991 | 108,000 | $\ldots$ | 1.88 | $\ldots$ | 0.025 | $\ldots$ |
| Hen $905 . . . . . . . . . . .$. | 1987 | 182,000 | $\ldots$ | 2.77 | ... | 0.025 | $\ldots$ |
| Hen 1103............ | 1987 | 183,000 | $\ldots$ | 2.54 | ... | 0.019 | $\ldots$ |
|  | 1991 | 131,000 | $\ldots$ | 2.80 | $\ldots$ | 0.049 | $\ldots$ |
| Hen 1213............ | 1987 | 112,000 | $\ldots$ | 1.84 | $\ldots$ | 0.022 | $\ldots$ |
|  | 1991 | 109,000 | $\ldots$ | 1.80 | ... | 0.022 | $\ldots$ |
|  | 1995 | 96,000 | $\ldots$ | 1.91 | $\ldots$ | 0.033 | $\ldots$ |
| KX $\operatorname{Tr} \mathrm{A} \ldots \ldots \ldots \ldots .$. | 1987 | 165,000 | $114^{\text {a }}$ | 3.04 | $3.9{ }^{\text {a }}$ | 0.041 | $\ldots$ |
|  |  |  | $140-150^{\text {b }}$ | ... | $2.8{ }^{\text {b }}$ | ... | $\ldots$ |
| AE Ara .............. | 1991 | 111,000 | 54-80 ${ }^{\text {a }}$ | 2.10 | $3.0{ }^{\text {a }}$ | 0.031 | $\ldots$ |
|  |  |  | $70-105^{\text {b }}$ | ... | $0.9{ }^{\text {b }}$ | ... | . |
| AS 255 | 1987 | 198,000 | ... | 3.04 | $\ldots$ | 0.028 | $\ldots$ |
|  | 1991 | 205,000 | $\ldots$ | 3.15 | .. | 0.030 | . |
| AS 270 | 1987 | 141,000 | ... | 1.76 | $\ldots$ | 0.013 | . |
|  | 1991 | 145,000 | $\ldots$ | 1.79 | $\ldots$ | 0.013 | $\ldots$ |
|  | 1995 | 124,000 | ... | 2.15 | $\ldots$ | 0.026 | $\ldots$ |
| AS 289 | 1987 | 174,000 | $\ldots$ | 0.90 | $\ldots$ | 0.003 | $\ldots$ |
|  | 1991 | 182,000 | ... | 1.23 | $\ldots$ | 0.004 | $\ldots$ |
| Y CrA ............... | 1990 | 190,000 | $130-140^{\text {b }}$ | 2.15 | $2.4{ }^{\text {b }}$ | 0.011 | ... |
| AS $304 \ldots \ldots . . . . . .$. | 1987 | 117,000 | $>65^{\text {c }}$ | 2.06 | $>3.9^{\text {c }}$ | 0.026 | $<0.68^{\text {c }}$ |
| AS 316 | 1987 | 159,000 | $114^{\text {a }}$ | 2.89 | $3.3{ }^{\text {a }}$ | 0.037 | ... |
|  | 1991 | 152,000 | $\ldots$ | 2.87 | $\ldots$ | 0.039 | $\ldots$ |
| AS 327 | 1987 | 192,000 | $114^{\text {a }}$ | 2.77 | $3.4{ }^{\text {a }}$ | 0.022 | $\ldots$ |
|  | 1991 | 177,000 | $\ldots$ | 2.82 | ... | 0.028 | $\ldots$ |
| CD $-43^{\circ} 14304 \ldots \ldots$ | 1987 | 167,000 | $\ldots$ | 2.38 | $\ldots$ | 0.019 | $\ldots$ |
|  | 1991 | 162,000 | $\ldots$ | 2.93 | $\ldots$ | 0.037 | $\ldots$ |

[^5]
### 6.3. Hen 1103

The behavior of the continuum in this object is similar to that of Hen 828, since it is higher in 1991 than in 1987. The dereddened spectra also behave like those of Hen 828, in the sense that in 1991 the continuum has an $\mathrm{A}-\mathrm{F}$ spectrum (Fig. 18). It is difficult to estimate its temperature, since the Balmer continuum is in emission, which does not allow a good judgment of the maximum intensity.

We consider that, as Hen 828, Hen 1103 was also having an outburst in 1991. There is, however, a difference between both objects, since the emission lines of Hen 1103 show no major changes between both dates of observation.

### 6.4. Hen 1213

This object was observed in the UV in 1995 (GutiérrezMoreno et al. 1997); the observations were severely affected by an extraordinarily high radiation level, but it was pos-
sible to state that lines of C IV, He II, and C III] were present.

### 6.5. KX TrA

According to Feibelman (1988) the low-dispersion International Ultraviolet Explorer (IUE) spectra of this object resemble strongly to V1016 Cyg and HM Sge; he states that the characteristics of the UV spectra suggest a highexcitation object containing a very hot star. The high excitation is also noticeable in the visual spectrum, rich in emission lines. There is a large difference between the reddening obtained by the Balmer decrement method and that obtained in UV, from the absence of the $\lambda 2200$ feature (Feibelman 1988); this low UV reddening is in agreement with the comment by Webster (1973) that the region is free of nebulosity and dust lanes and that any reddening is more likely due to circumstellar dust rather than interstellar dust. Determination of the reddening by using the He II lines in


Fig. 15.-H-R diagram of the objects observed. The curves are evolutionary tracks for stars of different masses (indicated at the right in units of $M_{\odot}$ ), as presented by Schönberner (1981; 0.546-0.644), Paczynski (1971; 0.8 and 1.2), and Shaw (1988; 1.4). Dots represent objects with $y \geq 0.17-$ which have been identified-whereas squares correspond to objects with $y<0.17$.
the visual ( $\lambda 4686$ ) and in $\mathrm{UV}(\lambda 1640)$ gives $\mathrm{C} \beta=0.48$, which is close to the mean between the visual and UV values of $\mathrm{C} \beta$.

### 6.6. AS 255

The dereddened spectra show that the continua corresponding to 1987 and 1991 have practically the same shape,


Fig. 16.-Dereddened spectra of Hen 828. The two spectra have the same zero point (separation step $=0$ ).


Fig. 17.-Tracing of the two-dimensional spectrum of Hen 905 through the lines O vi 26825 . It shows an emission within about $2^{\prime \prime}$ of the main object, with faint emissions in between.
but the 1987 spectrum is slightly brighter than that of 1991 . There is also a large difference in the intensity of the Balmer lines, which may be appreciated for $\mathrm{H} \beta$ in Tables 4 and 6.

### 6.7. AS 270

The dereddened spectra corresponding to 1987 and 1991 have practically identical continua, but that corresponding to 1995 shows an increase of the intensity shortward of 5000 $\AA$; at $4000 \AA$ the continuum intensity is about 3 times that of 1987. This fact suggests that the 1995 observations could also correspond to an outburst. Besides, the H lines are about twice as intense in the 1995 spectrum than in 1987 and 1991, whereas the observed fluxes of the other lines remain quite similar to those of the preceding epochs. This fact produces the decrease of the relative intensities of the lines of He I and He II, shown in Table 4, and the conse-


Fig. 18.-Same as Fig. 16, for Hen 1103
quent strong changes in the relative He abundance, shown in Table 7.

### 6.8. AS 289

The large intensity of Balmer lines, and mainly of $\mathrm{H} \beta$, in 1987, produces the difference in the relative He abundance, presented in Table 7.

According to Figure 15, the progenitor star of the hot component of this objects is the most massive of the sample.

### 6.9. AS 304

Munari \& Buson (1993) have analyzed IUE, optical, and infrared observations of this object. They find that the IUE spectra indicate that the hot component is losing mass at a high rate via a fast wind, being in this sense similar to the mass-losing nuclei of planetary nebulae ( PNe ).

They present a graph of the dereddened spectral energy distribution, which combines the $I U E$ continuum fluxes with $J H K L$ data, obtained almost simultaneously with the $I U E$ observations, and with the $U B V(R I)_{c}$ photometry obtained by Munari et al. (1992). The distribution obtained matches perfectly that of a bulge M4 giant longward of 7000 $\AA$, but departs from it at the $R$ band and shortward, due to the presence of the strong nebular continuum, which goes far into the red. The temperature, luminosity, and radius estimated by Munari \& Buson (1993) correspond to the He II Zanstra temperature.

### 6.10. AS 327

In this case, the variation in the He content reflects a real variation of the He lines, mainly $\lambda 5876$, and not a variation of the H lines intensity.

### 6.11. $\mathrm{CD}-43^{\circ} 14304$

The object was classified by Schmid \& Nussbaumer (1993) as a metal-poor yellow symbiotic. Schmid et al. (1998) have obtained a series of high-resolution optical spectra of this object, deriving the radial velocity curve of the cool component and an orbital period of about 1448 days. The hot component was also detected with ROSAT as a very soft X-ray source (Mürset, Wolff, \& Jordan 1997).

The dereddened spectra corresponding to 1991 June and July are slightly brighter than that corresponding to 1987; besides, both of them, show an increase of the intensity
shortward of $5000 \AA$, being the intensity at $4000 \AA$ in 1991 about 3 times that of 1987. This fact suggests that the objects was having a mild outburst in 1991.

## 7. GENERAL DISCUSSION

We presented here observations of 14 S-type symbiotic stars, obtaining some of their physical parameters. These parameters are full of uncertainties, which start with the observational errors of the measures-due mainly to the difficulty in the tracing of the continuum - and the many uncertainties already mentioned in the determination of the distances. For example, using Tsuji (1978) scale of bolometric magnitudes for M-type giants and the relevant data from Dyck, Lockwood, \& Capps (1974) instead of Frogel \& Whitford's (1987) scale, we get distances that are about twice as large as those listed in this paper; these errors have strong influence in the determination of the luminosities and, consequently, in the obtention of the radii.

Another source of errors is the adopted value of the reddening, with strong effects in many of the results obtained here. We must stress the large differences observed in some cases between optical and UV values, with the optical values being larger than those obtained with UV measures. This effect has been discussed elsewhere (Gutiérrez-Moreno et al. 1997), an in this paper it affects mainly KX TrA and AE Ara and may be some other objects for which no UV observations are available. It may be that, as Feibelman (1988) suggests, the difference is due to circumstellar reddening present in these objects; this would imply that this circumstellar reddening is not detected in the UV.

A further comment concerning the reddening determination is that reddenings obtained from the ratio $F(1640) /$ $F(4686)$ of He II have to be taken with care, since when the UV reddening obtained independently of the visual observations-as is the case when the $\lambda 2200$ feature is used-and the reddening in the optical region are different-as found here for KX TrA-the result obtained with the He II method is a combination of both values of $\mathrm{C} \beta$, and does not represent the real reddening.

Partial corrections for collisional effects were used for the determination of the He II abundances ( $\gamma \leq 0.3$, according to Peimbert \& Torres-Peimbert notation). The value of $\gamma$ was chosen so as to get the best agreement between the abundances obtained for the different lines used. Nevertheless, the corrections are very small when Almog \& Netzer (1989) emissivities are used; this indicates that their method takes very well into account the collisional corrections.

The H-R diagram presented in Figure 15 shows some evolutionary tracks for the hot components of central stars of PNe. Our objects are located in positions similar to those
of rather evolved PNe nuclei. They show a large scatter both in $T^{*}$ and $L^{*}$, with $96,000 K<T^{*}<205,000 \mathrm{~K}$ for Hen 1213 (1995 observation) and AS 255, respectively; and $1.0<\log L / L_{\odot}<3.2$, for AS 289 and AS 255, respectively. It is noticeable that the all objects with high He contentswith the exception of AE Ara-are located to the left of the evolutionary track corresponding to $M^{*}=0.644 M_{\odot}$, suggesting rather massive progenitors.

On the other hand, a look at the radii of the hot components given in Table 9, even considering possible errors of the order of a factor of 2 (due to the uncertainties of the distances), shows that they seem to be all white dwarfs.

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[^1]:    Note.-Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.
    ${ }^{\text {a }}$ Infrared values taken from Allen 1982a, 1982b.

[^2]:    ${ }^{2}$ IRAF is distributed by National Optical Astronomical Observatories, which is operated by the Association of Universities for Research in Astronomy, Inc., under contract to the National Science Foundation.

[^3]:    ${ }^{\text {a }}$ Badly blended with [O III] $\lambda 5007$.

[^4]:    ${ }^{\text {a }}$ Harries \& Howarth 1996
    ${ }^{\text {b }}$ Mikolajewska et al. 1997.
    c Pereira 1995.
    ${ }^{d}$ Munari \& Buson 1993.
    ${ }^{\text {e }}$ Schmid \& Nussbaumer 1993.

[^5]:    ${ }^{\text {a }}$ Mikolajewska et al. 1997.
    ${ }^{\text {b }}$ Pereira 1995.
    ${ }^{\text {c }}$ Munari \& Buson 1993.

