# Chemical Abundances of Main-sequence, Turnoff, Subgiant, and Red Giant Stars from APOGEE Spectra. I. Signatures of Diffusion in the Open Cluster M67 

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#### Abstract

Detailed chemical abundance distributions for 14 elements are derived for eight high-probability stellar members of the solar metallicity old open cluster M67 with an age of $\sim 4 \mathrm{Gyr}$. The eight stars consist of four pairs, with each pair occupying a distinct phase of stellar evolution: two G dwarfs, two turnoff stars, two G subgiants, and two red clump (RC) K giants. The abundance analysis uses near-IR high-resolution spectra ( $\lambda 1.5-1.7 \mu \mathrm{~m}$ ) from the Apache Point Observatory Galactic Evolution Experiment survey and derives abundances for C, N, O, Na, Mg, Al, Si, K, $\mathrm{Ca}, \mathrm{Ti}, \mathrm{V}, \mathrm{Cr}, \mathrm{Mn}$, and Fe . Our derived stellar parameters and metallicity for 2M08510076+1153115 suggest that this star is a solar twin, exhibiting abundance differences relative to the Sun of $\leqslant 0.04$ dex for all elements. Chemical homogeneity is found within each class of stars ( $\sim 0.02$ dex), while significant abundance variations ( $\sim 0.05-0.20$ dex) are found across the different evolutionary phases; the turnoff stars typically have the lowest abundances, while the RCs tend to have the largest. Non-LTE corrections to the LTE-derived abundances are unlikely to explain the differences. A detailed comparison of the derived $\mathrm{Fe}, \mathrm{Mg}, \mathrm{Si}$, and Ca abundances with recently published surface abundances from stellar models that include chemical diffusion provides a good match between the observed and predicted abundances as a function of stellar mass. Such agreement would indicate the detection of chemical diffusion processes in the stellar members of M67.


Key words: diffusion - infrared: stars - open clusters and associations: general - stars: abundances

## 1. Introduction

Messier 67 (M67, or NGC 2682) is one of the most-studied Galactic open clusters, due in part to it having an age and metallicity that are similar to those of the Sun, thus making it a useful cluster in which to study the properties of solar twins, as well as the evolution of solar-like stars. Photometric studies have provided a well-determined reddening $(E(B-V)=$ 0.041; Taylor 2007, Sarajedini et al. 2009), distance modulus ( $\mu=9.56-9.72$, or $d=800-860 \mathrm{pc}$; Yadav et al. 2008), and age ( $\sim 4.0$ Gyr; Salaris et al. 2004, Yadav et al. 2008, Sarajedini et al. 2009) for M67.

High-resolution spectroscopic abundance studies have found M67 to have a near-solar chemical composition ( $[\mathrm{Fe} / \mathrm{H}] \sim 0.00$; Cohen 1980; Foy \& Proust 1981; Friel \& Boesgaard 1992;

Tautvaišiene et al. 2000; Pancino et al. 2010; Jacobson et al. 2011; Önehag et al. 2014; Liu et al. 2016, Casamiquela et al. 2017). M67 also falls in one of the Kepler K2 campaign fields and thus has precision photometry available for a number of its members; these data have been used in recent asteroseismology studies of the red giants in M67 (e.g., Gonzalez 2016; Leiner et al. 2016; Stello et al. 2016).

In addition to containing stars that are similar to the Sun in mass, age, and chemical composition, M67 provides a laboratory in which to explore the properties of stars over a range of evolutionary phases: from dwarfs on the main sequence, up to the turnoff, through the subgiant branch (SGB), and onto the red giant branch (RGB; with these stars having very nearly the same ages and initial chemical

Table 1
Atmospheric Parameters

|  | 2M08510076 | 2M08512314 | 2M08514122 | 2M08505182 | 2M08513540 | 2M08514474 | 2M08521856 | 2M08514388 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | +1153115 | +1154049 | +1154290 | +1156559 | +1157564 | +1146460 | +1144263 | +1156425 |
|  | G Dwarf | G Dwarf | G Turnoff | G Turnoff | G Subgiant | G Subgiant | K Giant | K Giant |
| $B$ | 15.498 | 14.804 | 13.338 | 13.271 | 13.500 | 13.848 | 11.427 | 11.579 |
| V | 14.777 | 14.163 | 12.777 | 12.722 | 12.764 | 12.944 | 10.354 | 10.461 |
| $J$ | 13.474 | 13.017 | 11.703 | 11.646 | 11.427 | 11.357 | 8.572 | 8.618 |
| H | 13.157 | 12.741 | 11.466 | 11.382 | 11.143 | 10.918 | 8.087 | 8.114 |
| Ks | 13.105 | 12.681 | 11.397 | 11.321 | 11.030 | 10.822 | 7.923 | 7.996 |
| pm ra | $-12.5 \pm 2.5$ | $-9.5 \pm 1.6$ | $-8.7 \pm 1.3$ | $-11.7 \pm 1.3$ | $-9.7 \pm 1.3$ | $-10.3 \pm 1.2$ | $-11.9 \pm 1.5$ | $-9.7 \pm 2.5$ |
| pm dec | $-1.9 \pm 2.3$ | $-1.3 \pm 1.5$ | $-2.0 \pm 1.3$ | $-1.8 \pm 1.3$ | $-2.5 \pm 1.3$ | $-2.3 \pm 1.2$ | $-4.9 \pm 1.5$ | $-1.8 \pm 1.3$ |
| Probability | 98 | 99 | 100 | 100 | 91 | 99 | 96 | 97 |
| RV | $33.9 \pm 0.10$ | $33.6 \pm 0.10$ | $33.6 \pm 0.10$ | $30.8 \pm 0.10$ | $33.4 \pm 0.10$ | $33.1 \pm 0.10$ | $33.7 \pm 0.10$ | $32.9 \pm 0.1$ |
| SNR | 120 | 120 | 214 | 210 | 238 | 351 | 504 | 956 |
| $T_{\text {eff }}(\mathrm{K})$ | $5724 \pm 92$ | $5958 \pm 33$ | $6119 \pm 26$ | $6063 \pm 35$ | $5596 \pm 38$ | $5137 \pm 58$ | $4842 \pm 23$ | $4819 \pm 82$ |
| $\log g$ | $4.48 \pm 0.05$ | $4.35 \pm 0.05$ | $3.91 \pm 0.05$ | $3.87 \pm 0.05$ | $3.77 \pm 0.05$ | $3.64 \pm 0.05$ | $2.45 \pm 0.05$ | $2.44 \pm 0.05$ |
| $\xi \mathrm{km} \mathrm{s}^{-1}$ | 1.00 | 1.00 | 1.15 | 1.20 | 1.25 | 1.20 | 1.75 | 1.60 |

compositions). Quantitative high-resolution spectroscopic analyses can be used to reveal if there are chemical abundance differences between the M67 members in the different evolutionary phases, as well as determining whether these differences are real or if they reveal systematic differences induced by the analysis techniques themselves. Accurate chemical abundances of a large number of elements can be used to test models of chemical diffusion in stars (e.g., Michaud et al. 2015). The recent studies by Önehag et al. (2014), Blanco-Cuaresma et al. (2015), and Bertelli Motta et al. (2017) find chemical inhomogeneities for some elements that could be an indication that diffusion mechanisms may be at work and detectable in M67 stars.
Chemical abundance variations that result from diffusion have been probed in the metal-poor globular cluster NGC 6397 ( $[\mathrm{Fe} / \mathrm{H}]=-2.1$ ) in a series of papers by Korn et al. (2007), Lind et al. (2008), and Nordlander et al. (2012). These papers studied cluster stars at the turnoff point (TOP), on the SGB, at the base of the red giant branch (bRGB), and on the RGB using chemical abundances of $\mathrm{Li}, \mathrm{Mg}, \mathrm{Ca}, \mathrm{Ti}, \mathrm{Cr}$, and Fe . Chemical abundances were found to be dependent on evolutionary stage, and these were compared to diffusion models, with overall approximate agreement between observed and model abundances, particularly for Mg and Fe .

Diffusion model predictions specific to the solar metallicity and age regime of M67 have been presented by Michaud et al. (2004) and most recently by Dotter et al. (2017). In both studies, diffusion effects near the turnoff in M67 have been found to be as large as $\sim 0.1$ dex when compared to cooler main-sequence stars or stars on the RGB (where convection erases the abundance patterns created by diffusion), especially for certain elements such as Fe or Mg. Specific patterns in the abundance distributions of certain elements, such as Ca and Fe , are also predicted to exhibit detectable variations, suggesting that M67 is a key open cluster in which to probe and test models of diffusion processes.

In this study, a small sample of M67 members that were observed as part of the Apache Point Observatory Galactic Evolution Experiment (APOGEE; Majewski et al. 2017) survey are analyzed using its high-resolution, near-infrared (NIR) spectra to derive accurate chemical abundances for a large number of elements. The studied sample spans a range of evolutionary phases consisting of G dwarfs, G turnoff stars,

G subgiants, and K giants. The detailed chemical abundances derived for these stars are used to investigate possible abundance inhomogeneities in M67 as a function of stellar class and determine whether such abundance variations can be explained by a physical process, such as diffusion in the stellar envelope, or reflect systematic effects associated with the analysis techniques. The recent studies of Bovy (2016) and Price-Jones \& Bovy (2018) using APOGEE spectra found the tightest constraints on the chemical homogeneity of M67 red giants. This evidence of chemical homogeneity in red giants is an important starting point to investigate possible diffusion effects and assign observed abundance differences to the effects of stellar evolution.

Section 2 discusses the details of the APOGEE survey and spectra, while Section 3 presents the determination of stellar parameters (effective temperature, $T_{\text {eff }}$, and surface gravity, $\log g$ ), along with the chemical abundances. The chemical abundance distributions and possible variations are discussed in Sections 4 and 5, with the summary of results in Section 6.

## 2. The APOGEE Spectra and the Sample

The spectra analyzed in this work are from the SDSS-IV/ APOGEE2 survey (Blanton et al. 2017, Majewski et al. 2017). The APOGEE instrument is a cryogenic multifiber spectrograph (300 fibers; Wilson et al. 2010) on the SDSS 2.5 m telescope at the Apache Point Observatory (Gunn et al. 2006). The survey observations consist of high-resolution $(R=$ $\lambda / \Delta \lambda \sim 22,500)$ spectra of stars, primarily red giants but also stars in other evolutionary stages (see Zasowski et al. 2013), in the NIR $(\sim \lambda 1.50-\lambda 1.70 \mu \mathrm{~m})$ with the ultimate goal of exploring the chemical evolution of the stellar populations in the Milky Way.

Our sample contains eight stars that are deemed to be members of the M67 open cluster. The APOGEE spectra of the sample stars were reduced automatically by the APOGEE pipeline (Holtzman et al. 2015; Nidever et al. 2015) and then analyzed manually to extract detailed chemical abundances. We selected targets strategically to sample a range in effective temperatures and surface gravities that are representative of stars on the main sequence and in more advanced phases of evolution: two G dwarfs, two G-type turnoff stars, two G-type subgiants, and two K-type red giants. All targets are in the


Figure 1. 2MASS color-magnitude diagram of the APOGEE targets in the M67 field (shown as gray dots). The red symbols correspond to the studied stars: solar-like stars (red diamonds), turnoff stars (red squares), subgiant stars (red triangles), and red clump stars (red circles). Two isochrones for an age of $4 \mathrm{Gyr},(m-M)_{0}=9.60$, and $[\mathrm{Fe} / \mathrm{H}]=0.00$ from PARSEC (black line) and MIST (blue line) are also shown.
proper-motion study of Yadav et al. (2008) and have probabilities of membership higher than $91 \%$ (Table 1). Their measured radial velocities from the APOGEE spectra are also presented in Table 1, and these are consistent with probable cluster membership (Geller et al. 2015). Figure 1 shows the color-magnitude diagram for the NIR 2MASS data (Skrutskie et al. 2006), with $H_{0}$ plotted versus $\left(J-K_{s}\right)_{0}$. The eight target stars are shown as filled red symbols. We also show, as gray dots, the 536 stars from the M67 field observed by APOGEE (which include additional M67 members) and two isochrones for an age of $4 \mathrm{Gyr},(m-M)_{0}=9.60$, and $[\mathrm{Fe} / \mathrm{H}]=0.00$ from PARSEC (Bressan et al. 2012) and MIST (Choi et al. 2016; Dotter 2016).

## 3. Chemical Abundance Analysis

The sample studied here contains a mixture of stellar types. Both of the K giants in the sample were found to be red clump (RC) giants by Stello et al. (2016) based on $K 2$ asteroseismology. In addition, Stello et al. (2016) analyzed the $K 2$ oscillations from one of the stars, $2 \mathrm{M} 08511474+1146460$, and confirmed it to be a subgiant star very close to the base of the RGB.

The spectra of all eight stars were analyzed for chemical abundances in a homogeneous way that is independent from the methodology adopted in the derivation of stellar parameters and chemical abundances for the 14th SDSS Data Release (DR14) using the APOGEE automatic abundance pipeline ASPCAP (García Pérez et al. 2016). The "boutique" manual
abundance analysis described in this section provides independent results for a small sample that can be compared to the automated ASPCAP results for a larger sample of M67 members.

### 3.1. Effective Temperatures

The effective temperatures for the stars in the studied sample were derived using the photometric calibrations of González-Hernández \& Bonifacio (2009) and five different colors ( $B-V, V-J, V-H, V-K s$, and $J-K s$ ) while adopting a solar metallicity. The individual magnitudes $B$ and $V$ were taken from Yadav et al. (2008), and the infrared colors are from the 2MASS catalog (Skrutskie et al. 2006). A reddening $E(B-V)=0.041$ (Taylor 2007; Sarajedini et al. 2009) was adopted with individual dereddened colors obtained using the relations from Schlegel et al. (1998) and Carpenter (2001). The adopted colors and derived effective temperatures, plus the standard deviations of the mean (typically below $\sim 50 \mathrm{~K}$ ), are presented in Table 1. When the internal uncertainties in the González-Hernández \& Bonifacio (2009) calibration are included, along with the errors in the photometric colors, the total estimated uncertainty expected in $T_{\text {eff }}$ is $\sim 100 \mathrm{~K}$, adding all of the errors in quadrature.

### 3.2. Surface Gravities

Stellar surface gravities $(\log g)$ were derived from the fundamental relation with stellar mass, $T_{\text {eff }}$, and absolute
bolometric magnitude (Equation (1) below). Stellar masses were estimated by using the absolute magnitudes for the $B, V$, $H$, $J$, and $K s$ filters, along with an adopted PARSEC isochrone for an age $=4.0 \mathrm{Gyr},[\mathrm{Fe} / \mathrm{H}]=0.0$, and $(m-M)_{0}=9.60$, with derived masses then being $\sim 1.34 M_{\odot}$ for the RC, $\sim 1.30 M_{\odot}$ for the subgiants, $\sim 1.20 M_{\odot}$ for the turnoff stars, and $\sim 1.00 M_{\odot}$ for the solar-like stars. We note that if the isochrone from MIST (Figure 1) is adopted, the obtained stellar masses are not significantly different. The combination of effective temperature, stellar mass, and bolometric magnitudes (with bolometric corrections from Montegriffo et al. 1998) in Equation (1) provides the surface gravity values listed in Table 1. The solar values used were $\log g_{\odot}=4.438 \mathrm{dex}, T_{\text {eff, } \odot}=5772 \mathrm{~K}$, and $M_{\mathrm{bol}, \odot}=4.75$, which follows the IAU prescription in Prša et al. (2016):

$$
\begin{align*}
\log g= & \log g_{\odot}+\log \left(\frac{M_{\star}}{M_{\odot}}\right)+4 \log \left(\frac{T_{\star}}{T_{\odot}}\right) \\
& +0.4\left(M_{\mathrm{bol}, \star}-M_{\mathrm{bol}, \odot}\right) \tag{1}
\end{align*}
$$

Estimates of the uncertainties in the surface gravities were computed from two isochrones with ages of 3.5 and 4.5 Gyr that were used to rederive the stellar masses. Included in these estimates were the effective temperature uncertainties, as well as a typical metallicity uncertainty of $\pm 0.05$ dex. Errors in the photometric magnitudes are small and of the order of 0.03 mag. Combining all of these values in quadrature, an uncertainty of $\sim 0.07$ dex is found for the derived values of $\log g$. Our sample includes three stars with asteroseismic $\log g$ from the $K 2$ mission reported in Stello et al. (2016). Our derived $\log g$ values agree quite well with the asteroseismic ones; $\delta \log g$ (this work; Stello et al. 2016) $=0.04 \pm 0.07$ dex.

### 3.3. Chemical Abundances and Selected Lines

Chemical abundances were derived from a 1D LTE analysis and spectral synthesis using the Turbospectrum code (Alvarez \& Plez 1998; Plez 2012) in combination with model atmospheres interpolated from the MARCS ${ }^{21}$ grid (Gustafsson et al. 2008) for the atmospheric parameters derived in Sections 3.1 and 3.2 (Table 1). The APOGEE line list used in all computations was an updated version of the DR14 line list (line list 20170418). Shetrone et al. (2015) provided details on the construction of the APOGEE line list, and updates can be found in J. Holtzman et al. (2018, in preparation).

In a 1D abundance analysis, the microturbulent velocity $(\xi)$ is a necessary parameter to have both weak and strong lines of a given species yield the same abundance. Microturbulent velocities were determined using the same procedure as in Souto et al. (2016; or Smith et al. 2013). The determination of $\xi$ relies on Fe I lines that span a range of line strengths (or equivalent widths), with the stronger lines displaying a much larger sensitivity of the derived abundance with the microturbulent velocity. The best value of $\xi$ yields the closest agreement in the abundances of the strong and weak lines. In practice, adopted values of $\xi$ were varied from 0.5 to $3.0 \mathrm{~km} \mathrm{~s}^{-1}$. The inferred values of $\xi$ are included in Table 1.

Chemical abundances of individual elements were derived using a line-by-line manual analysis and obtaining best fits of synthetic spectra to observed line profiles. Local continuum levels in the observed spectra were set manually, and the particular

[^0]elemental abundance varied, until the differences between the observed and synthetic spectra were minimized. The instrumental resolution of the APOGEE spectrograph $(R \sim 22,500)$ produces an instrumental profile with an $\sim 13.7 \mathrm{~km} \mathrm{~s}^{-1}$ full-width at halfmaximum ( $\mathrm{FWHM} \sim 0.71 \AA$ ). Small variations in the broadening result in small adjustments across the spectra of about $\sim \pm 1.5 \mathrm{~km} \mathrm{~s}^{-1}$. During the line-profile fitting, searches were made for extra-broadening effects related to $v \sin (i)$ and/or macroturbulence; however, no extra line broadening was needed beyond the instrumental APOGEE profile to obtain good fits to the observed line profiles.

Chemical abundances of the elements $\mathrm{C}, \mathrm{N}, \mathrm{O}, \mathrm{Na}, \mathrm{Mg}, \mathrm{Al}, \mathrm{Si}$, $\mathrm{K}, \mathrm{Ca}, \mathrm{Ti}, \mathrm{V}, \mathrm{Cr}, \mathrm{Mn}$, and Fe were derived for the two RC giants in our sample (Table 2). The selected transitions were the same as in our previous study of red giants in the open cluster NGC 2420 using APOGEE spectra (Souto et al. 2016). Souto et al. (2016) analyzed a sample of red giants with similar values of $T_{\text {eff }}$ and $\log g$ (although slightly more metal-poor, with $[\mathrm{Fe} / \mathrm{H}]=-0.20$ ) to those of the M67 RC giants studied here.

The previous work on NGC 2420 by Souto et al. (2016) did not analyze solar-like dwarfs, turnoff stars, or subgiants. In this study, we have made a careful search for usable spectral lines or features in the APOGEE spectra of dwarfs and subgiants, with the goal being to maximize the number of lines available for each chemical species. Initial identification of promising spectral lines in solar-like dwarfs was done using an APOGEE spectrum of the asteroid Vesta (as a solar proxy), which were observed using the APOGEE spectrograph fiber-linked to the APO 1 m telescope.

The atomic and molecular lines used in the "boutique" abundance analysis of the stars (and their associated abundances) are listed in Table 2. A total of 135 spectral lines or features were selected as abundance indicators: $77 \mathrm{Fe} \mathrm{I}, 4 \mathrm{CO}$, $3 \mathrm{CI}, 10 \mathrm{CN}, 4 \mathrm{OH}, 2 \mathrm{NaI}, 6 \mathrm{Mg}$, $3 \mathrm{AlI}, 9 \mathrm{SiI}, 2 \mathrm{KI}, 4 \mathrm{CaI}$, 6 Ti I, 1 VI , 1 CrI , and 3 Mn . As the studied sample covers an extended range in $T_{\text {eff }}-\log g$ parameter space, it is not possible to measure the exact same transitions for all stars, given that the strengths of the various spectral features change as a function of $T_{\text {eff }}$ and $\log g$.

The APOGEE spectra of cool red giants are typically dominated by molecular features (mostly $\mathrm{CO}, \mathrm{CN}$, and OH ) but also show atomic lines from many elements. For G dwarfs and subgiants with higher effective temperatures, molecular absorption becomes less important, and neutral atomic lines become dominant. In the case of G dwarfs and turnoff stars, it is not possible to derive oxygen and nitrogen abundances, as the molecular OH and CN lines become too weak. Vanadium abundances are also not measurable. Most chemical abundances from APOGEE are based on transitions of neutral atomic lines and include $\mathrm{Fe}, \mathrm{Na}, \mathrm{Al}, \mathrm{Mg}, \mathrm{Si}, \mathrm{K}, \mathrm{Ca}, \mathrm{Ti}, \mathrm{V}, \mathrm{Cr}$, and Mn . These lines can be measured, for the most part, in all four stellar classes, although for some lines, the degree of blending and/or the line strengths change significantly depending on the stellar class. Concerning iron, often taken as a metallicity indicator, a relatively small number of clean Fe I lines can be analyzed in red giants, while more than 70 Fe I lines can be used in the warmer subgiants, turnoff stars, and G dwarfs. Although the CO lines become quite weak for solartype stars, three C I lines ( $\lambda 15784.7$, $\lambda 16005.0$, and $\lambda 16021.7$ ) become stronger and measurable in these stars.

Figure 2 illustrates the observed and best-fit synthetic spectra for one star in each class: one RC, one subgiant, one turnoff star, and one solar-like G dwarf (top to bottom). These spectra

Table 2
Individual Abundances



| Element | $\lambda(\AA)$ | $\begin{gathered} \text { 2M08510076 } \\ +1153115 \\ \text { G Dwarf } \end{gathered}$ | $\begin{gathered} \text { 2M08512314 } \\ \text { +1154049 } \\ \text { G Dwarf } \end{gathered}$ | $\begin{gathered} \text { 2M08514122 } \\ +1154290 \\ \text { G Turnoff } \end{gathered}$ | $\begin{gathered} \text { 2M08505182 } \\ +1156559 \\ \text { G Turnoff } \end{gathered}$ | $\begin{gathered} \text { 2M08513540 } \\ +1157564 \\ \text { G Subgiant } \end{gathered}$ | $\begin{gathered} \text { 2M08514474 } \\ +1146460 \\ \text { G Subgiant } \end{gathered}$ | $\begin{gathered} \text { 2M08521856 } \\ +1144263 \\ \text { K Giant } \end{gathered}$ | $\begin{gathered} \text { 2M08514388 } \\ +1156425 \\ \text { K Giant } \end{gathered}$ | Sun <br> This Work |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 15447. | $\ldots$ | $\ldots$ | $\ldots$ | $\cdots$ | 8.07 | 7.93 | 8.17 | 8.12 | $\cdots$ |
|  | 15466. | ... | ... | ... | ... | 8.03 | 7.95 | 8.16 | 8.06 | ... |
|  | 15472. | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | 7.90 | 8.18 | 8.07 | ... |
|  | 15482. | ... | $\ldots$ | ... | ... | ... | 7.96 | 8.17 | 8.06 | $\ldots$ |
|  | 15580.88 | ... | ... | ... | ... | 8.09 | 8.05 | 8.18 | 8.09 | ... |
| OH | 15278.334 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 8.60 | 8.63 | $\cdots$ | $\cdots$ | $\cdots$ |
|  | 15568.780 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 8.66 | 8.68 | $\ldots$ | ... | $\ldots$ |
|  | 16190.263 | $\ldots$ | ... | $\ldots$ | ... | 8.68 | 8.65 | $\ldots$ | $\ldots$ | ... |
|  | 16192.208 | ... | ... | ... | ... | 8.63 | 8.65 | ... | ... | ... |
| Na I | 16373.853 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 6.37 | 6.40 | 6.40 | 6.37 |  |
|  | 16388.858 | 6.33 | ... | 6.29 | 6.36 | 6.38 | 6.42 | 6.41 | 6.46 | 6.31 |
| Mg I | 15740.716 | 7.43 | 7.48 | 7.45 | 7.33 | 7.49 | 7.45 | 7.58 | 7.60 | 7.45 |
|  | 15748.988 | 7.44 | 7.48 | 7.38 | 7.34 | 7.50 | 7.48 | 7.65 | 7.62 | 7.44 |
|  | 15765.842 | 7.41 | 7.46 | 7.36 | 7.35 | 7.52 | 7.48 | 7.54 | 7.52 | 7.46 |
|  | 15879.5 | $\ldots$ | ... | ... | ... | $\ldots$ | ... | 7.50 | 7.57 | $\ldots$ |
|  | 15886.2 | $\ldots$ | ... | ... | $\ldots$ | $\ldots$ | ... | 7.72 | 7.71 | $\ldots$ |
|  | 15954.477 | ... | ... | ... | ... | ... | ... | 7.68 | 7.62 | ... |
| Al I |  |  |  |  |  |  |  |  | $6.56$ | $6.36$ |
|  | $16750.564$ | 6.37 | 6.33 | 6.32 | 6.38 | 6.43 | 6.38 | 6.57 | 6.55 | 6.34 |
|  | 16763.360 | 6.44 | 6.47 | 6.42 | 6.37 | 6.51 | 6.46 | 6.59 | ... | 6.42 |
| Si I | 15361.161 | $\ldots$ | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 7.62 | 7.63 | $\ldots$ |
|  | 15376.831 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | 7.61 | 7.59 | $\ldots$ |
|  | 15888.410 | 7.46 | 7.52 | 7.46 | 7.36 | 7.52 | 7.48 | $\ldots$ | ... | 7.42 |
|  | 15960.063 | 7.43 | 7.45 | 7.45 | 7.39 | 7.53 | 7.52 | ... | $\ldots$ | 7.46 |
|  | 16060.009 | ... | ... | $\ldots$ | ... | ... | ... | 7.56 | 7.59 | ... |
|  | 16094.787 | 7.45 | 7.44 | 7.46 | 7.39 | 7.49 | 7.48 | 7.65 | 7.63 | 7.45 |
|  | 16215.67 | . | $\ldots$ | ... | ... | $\ldots$ | $\ldots$ | 7.63 | 7.61 | $\ldots$ |
|  | $16680.770$ | 7.50 | 7.54 | 7.45 | 7.52 | 7.54 | 7.49 | 7.67 | $7.64$ | 7.53 |
|  | $16828.159$ | . | . | . | . | . | . | 7.60 | 7.62 | ... |
| K I | 15163.067 | 5.08 | 5.09 | 5.07 | 5.17 | 5.12 | 5.08 | 5.14 | 5.04 | 5.06 |
|  | 15168.376 | 5.13 | 5.16 | 5.14 | 5.08 | 5.11 | 5.06 | 5.13 | 5.05 | 5.10 |
| Ca I | $16136.823$ | 6.38 | 6.38 | 6.16 | 6.15 |  | 6.32 | 6.33 | 6.33 | 6.36 |
|  | $16150.763$ | 6.32 | 6.36 | ... | ... | 6.30 | 6.33 | 6.37 | 6.36 | 6.31 |
|  | 16155.236 | ... | .. | $\ldots$ | $\ldots$ | ... | $\ldots$ | 6.43 | 6.39 | ... |
|  | 16157.364 | 6.37 | 6.33 | 6.26 | 6.20 | 6.36 | 6.33 | 6.39 | 6.40 | 6.32 |
| Ti I | 15334.847 | .. | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
|  | 15543.756 | 4.92 | ... | ... | ... | 4.86 | 4.87 | 4.97 | 4.91 | 4.92 |
|  | 15602.842 | .. | $\ldots$ | $\cdots$ | $\cdots$ | $\cdots$ | 4.89 | 5.04 | 4.99 | ... |




Figure 2. Best-fit synthetic spectra (black) overplotted with the observed spectra (blue) for four target stars: one red clump, one subgiant, one turnoff, and one solar-like star (top to bottom). The spectral lines used to derive the abundances are marked with black circles.
highlight the region between 16140 and 16270 Å, covering only a small portion of the APOGEE wavelength range. Many of the spectral lines in Figure 2 become noticeably broader in the G-dwarf spectra due primarily to increasing surface gravity. The nature of the absorption lines changes noticeably when going from the K giant (top panel) down to the G dwarf (bottom panel) with the weakening of the molecular lines.

Abundance results from line-by-line measurements are shown in Table 2, while mean elemental abundances obtained for each star are presented in Table 3.

### 3.4. Abundance Uncertainties

To estimate the uncertainties in the derived abundances due to the uncertainties in the adopted stellar parameters, new abundances were computed for perturbed values of the microturbulent velocity, as well as using model atmospheres with perturbed values of effective temperature, surface gravity, and metallicity.

For all stellar classes, the baseline models corresponded to the stars: $2 \mathrm{M} 08510076+1153115$ (solar-like), 2M08505182 +1156559 (turnoff), $2 \mathrm{M} 08513540+1157564$ (subgiant), and

Table 3
Stellar Abundances

|  | 2M08510076 | 2M08512314 | 2M08514122 | 2M08505182 | 2M08513540 | 2M08514474 | 2M08521856 | 2M08514388 | Sun |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | +1153115 | +1154049 | +1154290 | +1156559 | +1157564 | +1146460 | +1144263 | +1156425 | This Work |
|  | G Dwarf | G Dwarf | G Turnoff | G Turnoff | G Subgiant | G Subgiant | K Giant | K Giant |  |
| Fe | $7.46 \pm 0.04$ | $7.49 \pm 0.04$ | $7.40 \pm 0.04$ | $7.38 \pm 0.03$ | $7.47 \pm 0.03$ | $7.44 \pm 0.03$ | $7.50 \pm 0.03$ | $7.49 \pm 0.03$ | $7.45 \pm 0.03$ |
| C | $8.41 \pm 0.02$ | $8.39 \pm 0.02$ | $8.28 \pm 0.02$ | $8.33 \pm 0.04$ | $8.35 \pm 0.05$ | $8.38 \pm 0.07$ | $8.34 \pm 0.01$ | $8.37 \pm 0.02$ | $8.37 \pm 0.05$ |
| N | ... | ... | ... | ... | $8.08 \pm 0.04$ | $7.96 \pm 0.05$ | $8.15 \pm 0.04$ | $8.07 \pm 0.02$ | ... |
| O | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | ... | $8.64 \pm 0.03$ | $8.65 \pm 0.02$ | $\ldots$ |
| Na | 6.33 | $\cdots$ | 6.29 | 6.36 | $6.38 \pm 0.01$ | $6.41 \pm 0.01$ | $6.41 \pm 0.01$ | $6.41 \pm 0.07$ | 6.31 |
| Mg | $7.43 \pm 0.01$ | $7.48 \pm 0.01$ | $7.40 \pm 0.04$ | $7.34 \pm 0.01$ | $7.50 \pm 0.02$ | $7.47 \pm 0.01$ | $7.61 \pm 0.07$ | $7.61 \pm 0.06$ | $7.44 \pm 0.01$ |
| Al | $6.40 \pm 0.03$ | $6.39 \pm 0.06$ | $6.38 \pm 0.04$ | $6.37 \pm 0.01$ | $6.44 \pm 0.05$ | $6.42 \pm 0.03$ | $6.57 \pm 0.02$ | $6.55 \pm 0.01$ | $6.37 \pm 0.04$ |
| Si | $7.46 \pm 0.02$ | $7.51 \pm 0.03$ | $7.46 \pm 0.01$ | $7.42 \pm 0.06$ | $7.52 \pm 0.02$ | $7.50 \pm 0.01$ | $7.62 \pm 0.03$ | $7.62 \pm 0.02$ | $7.47 \pm 0.04$ |
| K | $5.11 \pm 0.03$ | $5.13 \pm 0.04$ | $5.11 \pm 0.04$ | $5.18 \pm 0.05$ | $5.12 \pm 0.01$ | $5.07 \pm 0.01$ | $5.14 \pm 0.01$ | $5.05 \pm 0.01$ | $5.08 \pm 0.02$ |
| Ca | $6.37 \pm 0.03$ | $6.37 \pm 0.02$ | $6.21 \pm 0.05$ | $6.18 \pm 0.03$ | $6.34 \pm 0.03$ | $6.33 \pm 0.01$ | $6.38 \pm 0.04$ | $6.37 \pm 0.03$ | $6.33 \pm 0.03$ |
| Ti | $4.90 \pm 0.02$ | 4.93 | 4.96 | ... | $4.86 \pm 0.01$ | $4.88 \pm 0.04$ | $4.99 \pm 0.03$ | $4.94 \pm 0.04$ | 4.92 |
| V | ... | ... | ... | $\ldots$ | ... | ... | 4.12 | 4.06 | ... |
| Cr | 5.69 | 5.67 | 5.71 | 5.64 | 5.67 | 5.62 | 5.68 | 5.67 | 5.67 |
| Mn | $5.40 \pm 0.02$ | $5.39 \pm 0.02$ | $5.28 \pm 0.04$ | $5.26 \pm 0.03$ | $5.40 \pm 0.01$ | $5.38 \pm 0.02$ | $5.53 \pm 0.01$ | $5.51 \pm 0.03$ | $5.40 \pm 0.01$ |

$2 \mathrm{M} 08521856+1144263$ (RC). In each case, $T_{\text {eff }}$ was changed by $+100 \mathrm{~K}, \log g_{1}$ by +0.20 dex, $[\mathrm{Fe} / \mathrm{H}]$ by +0.20 dex, and $\xi$ by $+0.20 \mathrm{~km} \mathrm{~s}^{-1}$. Table 4 presents the final estimated abundance uncertainties, $\sigma$, representative of each stellar class. The final uncertainties were computed from the sum in quadrature of all the estimated uncertainties (same procedure as Souto et al. 2016, 2017).

The final abundance uncertainties for all stars are overall similar. The abundances of $\mathrm{Mg}, \mathrm{Al}$, and Si are more sensitive to changes in both $T_{\text {eff }}$ and $\log g$; their uncertainties are $\sim 0.10$ dex. In solar-like stars, the most sensitive elements to change in atmospheric parameters are Mg and Al , while for subgiants, the abundances of Mg and Si exhibit higher sensitivity to changes in the atmospheric parameters. The abundances of RC indicate a larger dependence on changes in $\xi$ for the Fe I and OH lines. Titanium is found to have the highest sensitivity to $T_{\text {eff }}$, while the oxygen abundance from the OH lines is found to be more sensitive to changes in $[\mathrm{Fe} / \mathrm{H}]$. Variations of $\log g$ by +0.20 dex in the model do not change significantly with the abundances in red giants.

### 3.5. Comparisons with Optical Studies from the Literature

Figures 3 and 4 show comparisons of the abundances obtained for all elements studied here with results from M67 high-resolution optical studies in the literature by Önehag et al. (2014; Liu et al. (2016; solar twins) and solar-like, turnoff, and subgiant stars) and the studies of red giants by Tautvaišiene et al. (2000), Yong et al. (2005), Friel et al. (2010), and Pancino et al. (2010). All abundances are plotted as a function of $\log g$ (Figure 3) and $T_{\text {eff }}$ (Figure 4), with the $[\mathrm{X} / \mathrm{H}]$ values from this study using our solar results as a reference, while $[\mathrm{X} / \mathrm{H}]$ values from the literature were taken directly from the other studies using their solar abundances as references. We note that these literature studies have analyzed the solar spectrum themselves (except for Önehag et al. 2014), and they use their derived solar abundances to measure differential abundances relative to the Sun. Our solar abundance results are also presented in Tables 2 and 3 .

Inspection of Figures 3 and 4 shows that our abundances for solar-like stars in M67 are in overall good agreement with those from both Liu et al. (2016) and Önehag et al. (2014). For the turnoff stars, we obtain results that are also similar to the ones
from Önehag et al. (2014); our abundances of $\mathrm{Mg}, \mathrm{Ca}$, and Mn are systematically lower by roughly $0.05-0.10$ dex. For the subgiant stars, our results agree in general with those of Önehag et al. (2014); the largest offset is found for Na , where we obtain an average Na abundance 0.09 dex higher than that in Önehag et al. (2014).

Several studies have determined chemical abundances of red giants in M67 via high-resolution spectroscopy, and a few of these works are compared in Figures 3 and 4. Overall, our RC abundances for all elements are similar to the literature values, falling close to the middle of the distribution of the abundances for the RC in Figures 3 and 4, except for the elements N, Al, Ca , and Mn (and to a certain degree Ti), for which our results fall in the upper envelope of the distribution by $\sim 0.05-0.20$ dex. We also note significant abundance scatter in the optical literature results for $\mathrm{O}, \mathrm{Na}, \mathrm{Ca}, \mathrm{Ti}, \mathrm{V}$, and Mn .

It is also noted that the abundances of $\mathrm{Mg}, \mathrm{Al}$, and Si in red giants are found to be, on average, enhanced relative to solar. Concerning metallicities, the mean iron abundance obtained for the two RC stars studied here is $\langle[\mathrm{Fe} / \mathrm{H}]\rangle=+0.05$, and these results agree with those of Pancino et al. (2010) but are slightly higher than the average metallicity from Tautvaišiene et al. (2000; $\langle[\mathrm{Fe} / \mathrm{H}]\rangle=-0.03$ ). Overall, the abundance patterns derived from the literature studies are similar to those obtained from the APOGEE spectra (with some significant offsets or scatter in certain elements). It is worth noting here again that all studies, this work included, are based on 1D LTE analyses.

## 4. Abundances Trends

Chemical abundance trends may result from a combination of simplification in the analysis or possible physical effects, such as non-LTE, 3D, and/or diffusion, as well as systematic errors in the abundance analysis. This subsection will discuss observed abundance trends across stars in different evolutionary phases for the abundance results presented in this study. This discussion will not include C and N , as abundance changes due to the first dredge-up (FDU) dominate any other physical processes, such as diffusion (Section 5.1). In addition, O and V are not included, as abundances from these elements are only measurable in the APOGEE spectra for red giants. This leaves 10 elements to investigate as a function of evolutionary state; simply put, stellar evolution proceeds

Table 4
Abundance Sensitivities Due to Atmospheric Parameters

| Stellar Class | Element | $\begin{gathered} T_{\text {eff }} \\ (+100 \mathrm{~K}) \end{gathered}$ | $\begin{gathered} \log g \\ (+0.2 \text { dex }) \end{gathered}$ | $\begin{gathered} \xi \\ \left(+0.2 \mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} {[\mathrm{M} / \mathrm{H}]} \\ (+0.2 \mathrm{dex}) \end{gathered}$ | $\sigma$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Red Giants | C | +0.03 | +0.02 | -0.03 | +0.04 | 0.062 |
|  | N | -0.04 | +0.02 | +0.00 | +0.08 | 0.092 |
|  | O | +0.05 | -0.03 | -0.06 | +0.11 | 0.138 |
|  | Na | +0.03 | -0.02 | +0.00 | +0.02 | 0.041 |
|  | Mg | +0.04 | -0.02 | $+0.00$ | +0.04 | 0.060 |
|  | Al | +0.08 | -0.02 | -0.04 | +0.04 | 0.100 |
|  | Si | +0.03 | -0.01 | -0.02 | +0.05 | 0.062 |
|  | K | +0.04 | -0.04 | -0.02 | +0.01 | 0.061 |
|  | Ca | +0.05 | -0.02 | -0.02 | +0.02 | 0.061 |
|  | Ti | +0.13 | +0.00 | -0.01 | +0.05 | 0.140 |
|  | V | +0.06 | +0.00 | -0.03 | +0.03 | 0.073 |
|  | Cr | +0.05 | -0.02 | -0.03 | +0.02 | 0.065 |
|  | Mn | +0.03 | +0.02 | -0.01 | +0.01 | 0.039 |
|  | Fe | +0.03 | -0.02 | -0.05 | +0.03 | 0.069 |
| Subgiants | C | +0.02 | $+0.02$ | -0.02 | $+0.00$ | 0.035 |
|  | N | -0.02 | +0.05 | -0.03 | +0.01 | 0.062 |
|  | O | +0.01 | -0.03 | -0.02 | +0.03 | 0.048 |
|  | Na | +0.01 | +0.01 | +0.00 | +0.01 | 0.017 |
|  | Mg | -0.06 | -0.07 | -0.02 | -0.09 | 0.130 |
|  | Al | -0.05 | -0.03 | -0.02 | -0.05 | 0.079 |
|  | Si | -0.05 | -0.03 | -0.03 | -0.04 | 0.155 |
|  | K | +0.02 | +0.01 | +0.00 | +0.00 | 0.077 |
|  | Ca | +0.01 | +0.01 | -0.01 | +0.00 | 0.022 |
|  | Ti | -0.02 | -0.02 | -0.03 | -0.03 | 0.017 |
|  | V | +0.03 | +0.00 | -0.01 | +0.00 | 0.014 |
|  | Cr | +0.02 | +0.02 | +0.00 | +0.00 | 0.028 |
|  | Mn | +0.02 | +0.01 | -0.01 | +0.00 | 0.024 |
|  | Fe | +0.02 | +0.07 | -0.03 | +0.01 | 0.079 |
| Turnoff | C | $+0.00$ | $+0.02$ | $+0.00$ | $+0.01$ | 0.022 |
|  | N | ... | ... | ... | ... | ... |
|  | O | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ |
|  | Na | +0.01 | +0.01 | $+0.00$ | +0.01 | 0.017 |
|  | Mg | +0.04 | -0.04 | +0.02 | -0.04 | 0.099 |
|  | Al | +0.02 | -0.03 | +0.02 | -0.05 | 0.076 |
|  | Si | +0.03 | -0.03 | +0.01 | -0.03 | 0.055 |
|  | K | +0.01 | +0.01 | +0.00 | +0.01 | 0.017 |
|  | Ca | +0.02 | +0.03 | +0.00 | +0.03 | 0.041 |
|  | Ti | +0.02 | -0.03 | +0.02 | -0.02 | 0.046 |
|  | V | ... | $\cdots$ | $\cdots$ | ... | ... |
|  | Cr | +0.00 | +0.01 | +0.00 | +0.01 | 0.014 |
|  | Mn | +0.02 | +0.00 | +0.01 | +0.01 | 0.024 |
|  | Fe | +0.03 | -0.02 | +0.02 | -0.02 | 0.046 |
| Solar-like | C | $+0.00$ | $+0.02$ | $+0.00$ | $+0.01$ | 0.022 |
|  | N | $\ldots$ | ... | ... | ... | ... |
|  | O | $\ldots$ | . | ... | ... | $\ldots$ |
|  | Na | +0.01 | +0.01 | +0.00 | +0.01 | 0.017 |
|  | Mg | +0.03 | -0.07 | +0.02 | -0.06 | 0.099 |
|  | Al | +0.02 | -0.05 | +0.02 | -0.05 | 0.076 |
|  | Si | +0.02 | -0.04 | +0.01 | -0.03 | 0.055 |
|  | K | +0.01 | +0.01 | +0.00 | +0.01 | 0.017 |
|  | Ca | +0.02 | +0.02 | +0.00 | +0.03 | 0.041 |
|  | Ti | +0.02 | -0.03 | +0.02 | -0.02 | 0.046 |
|  | V | $\cdots$ | $\cdots$ | $\cdots$ | $\cdots$ | $\ldots$ |
|  | Cr | +0.00 | +0.01 | +0.00 | +0.01 | 0.014 |
|  | Mn | +0.02 | +0.00 | +0.01 | +0.01 | 0.024 |
|  | Fe | +0.02 | -0.03 | +0.02 | -0.02 | 0.046 |

monotonically in $\log g$ from high $\log g$ to decreasing $\log g$ as the star evolves to the RGB phase.

Focusing on the abundance results obtained here as a function of $\log g$ (red filled symbols in Figure 3), two main
types of behavior can be noted: in one case, there is almost no change in the abundances as a function of surface gravity, which occurs for $\mathrm{K}, \mathrm{Ti}$, and Cr . In the other case, there are general increases in the abundances when comparing the


Figure 3. Chemical abundances for the studied stars as a function of $\log g$. Results from the literature for other stars in M67 from Liu et al. (Tautvaišiene et al. (2000) Yong et al. (2005), Friel et al. (2010), Pancino et al. (2010), 2016), and Önehag et al. (2014), are also show for comparison. The results from this study are shown as filled red symbols as in Figure 1.
main-sequence, turnoff, and subgiant stars with the RC stars (decreasing $\log g$ ). Among these seven elements, five have larger trends than the other two when comparing the less
evolved stars and the RC stars. These larger trends with surface gravity ( $\sim 0.20$ dex) are found for the abundances of the elements $\mathrm{Mg}, \mathrm{Al}, \mathrm{Si}, \mathrm{Ca}$, and Mn , with the mean abundances of


Figure 4. Same as Figure 3 but as a function of $T_{\text {eff. }}$.
the turnoff and RC stars being $\langle A(\mathrm{Mg})\rangle=7.37$ and $\langle A(\mathrm{Mg})\rangle=$ $7.61,\langle A(\mathrm{Al})\rangle=6.38$ and $\langle A(\mathrm{Al})\rangle=6.56,\langle A(\mathrm{Si})\rangle=7.44$ and $\langle A(\mathrm{Si})\rangle=7.62,\langle A(\mathrm{Ca})\rangle=6.20$ and $\langle A(\mathrm{Ca})\rangle=6.38$, and $\langle A(\mathrm{Mn})\rangle=5.27$ and $\langle A(\mathrm{Mn})\rangle=5.52$ for turnoff and red clump stars, respectively. The abundances of Na and Fe also
show a trend with $\log g$ but with a smaller change ( $\delta \sim 0.1$ dex $)$. For sodium and iron, we obtain $\langle A(\mathrm{Na})\rangle=$ 6.33 and $\langle A(\mathrm{Na})\rangle=6.41 \quad$ and $\langle A(\mathrm{Fe})\rangle=7.39 \quad$ and $\langle A(\mathrm{Fe})\rangle=7.50$ for the turnoff and RC, respectively. Three of these elements, $\mathrm{Na}, \mathrm{Al}$, and Si , display the smallest differences
between the abundances of the turnoff and solar-like stars and show a more monotonic increase with decreasing $\log g$.

Such trends in the abundances as a function of $\log g$ are also found when examining the literature optical results (shown in Figure 3), indicating behavior reminiscent of that found here for the APOGEE results. For example, Liu et al. (2016) studied solar twins in M67 and derived $[\mathrm{Na} / \mathrm{H}]=-0.05 \pm 0.02$, while the study of red giants by Friel et al. (2010) obtained $[\mathrm{Na} / \mathrm{H}]=+0.11 \pm 0.10$. The Al abundance for the solar twins from Liu et al. (2016) is $[\mathrm{Al} / \mathrm{H}]=+0.01 \pm 0.03$, while for the red giants Friel et al. (2010) obtained $[\mathrm{Al} / \mathrm{H}]=$ $+0.08 \pm 0.07$. Taking the Mg abundances as a second example and the literature results from Önehag et al. (2014) for the turnoff stars and Pancino et al. (2010) for the red giants, we see a change in the abundances of 0.31 dex in the mean.

Overall, the abundance results for most of the elements from all studies, taken at face value, would suggest that the abundances may increase as the star evolves (not expected), or, alternatively, that other effects become significant in the red giant regime. However, it must also be kept in mind that comparisons of results from different studies using different techniques and analysis methods may result in systematic differences in the abundances.

The distribution of the elemental abundances as a function of the effective temperature (Figure 4) can also be divided in elements that show small-to-no trends with $T_{\text {eff }}$ (such as $\mathrm{K}, \mathrm{Ti}$, and Cr ) and those that show a trend of increasing abundances with $T_{\text {eff }}$ between the solar-like and the K-type red giants (which may be related to surface gravity): $\mathrm{Mg}, \mathrm{Al}, \mathrm{Si}$, and Mn with the largest trends, while Na and Fe display a smaller effect. Calcium again shows a change in abundance that is driven exclusively by the low abundances in the turnoff stars. The abundances of $\mathrm{Mg}, \mathrm{Ca}, \mathrm{Mn}, \mathrm{Fe}$, and, to a lesser degree, Si show a pronounced decrease in the abundances of the turnoff stars (higher $T_{\text {eff }}$ ) when compared to the abundances of the solar-like and subgiant stars.

### 4.1. Departures from LTE

The trends of elemental abundances with surface gravity and effective temperature seen for some elements in Figures 3 and 4 , respectively, can be due to a combination of effects that include, but are not limited to, abundance offsets due to departures from LTE. This is a possible effect for the abundances in this study, as the targets cover a large range in $T_{\text {eff }}-\log g$ parameter space, although M67 has solar metallicity and departures from LTE are expected to be more significant for red giants in the metal-poor regime (e.g., Asplund 2005; Asplund et al. 2009). Several studies in the literature have investigated non-LTE effects for lines in the optical (e.g., Korn et al. 2007; Andrievsky et al. 2008; Lind et al. 2011; Bergemann et al. 2012; Osorio \& Barklem 2016; Smiljanic et al. 2016); however, to date, few non-LTE studies have investigated the behavior of transitions in the $H$ band and, in particular, in the APOGEE region.

Cunha et al. (2015) presented non-LTE abundance corrections for Na I lines in the APOGEE region for stars in the RC and on the RGB in the very metal-rich $([\mathrm{Fe} / \mathrm{H}]=+0.35)$ open cluster NGC 6791; the departures from LTE were found to be minimal. In this paper, the differences in $A(\mathrm{Na})$ between the red giants (RG) and the solar-like (SL) are $\Delta(\mathrm{RG}-\mathrm{SL})=$ +0.08 dex. Non-LTE effects are expected to be smaller than
these differences and of the order of $\sim 0.02$ dex (Lind et al. 2011; private communication).

Two recent studies investigated the formation of Mg I lines (Zhang et al. 2017) and Si I lines (Zhang et al. 2016) in the APOGEE region. Zhang et al. (2017) found that the Mg I lines $\lambda 15740, \lambda 15748$, and $\lambda 15765$ are well-modeled in LTE (showing only small non-LTE departures). We find that the Mg abundances in the RC stars are 0.24 dex larger than those in the turnoff stars. Such offsets between the abundances are unlikely to be explained as due to departures from LTE.

The results in Zhang et al. (2016) indicated that the Si I lines analyzed here at $15888,16380,16680$, and $16828 \AA$ show departures from LTE: $\delta A(\mathrm{Si})$ (non-LTE-LTE) $=-0.06$ for K-type red giant stars and $\delta A(\mathrm{Si})($ non-LTE-LTE $)=-0.05$ and -0.03 for G-type subgiants and solar-like stars, respectively. These non-LTE corrections are all in the same sense (all negative), and the offsets are roughly the same for red giants and subgiants and slightly smaller (by 0.03 dex) for solar-like stars. These non-LTE corrections of $\sim 0.01-0.03$ dex, although reducing the discrepancy, would not seem to be a plausible explanation for the 0.18 dex differences found here between the Si abundances of the red giants and turnoff stars.

## 5. Discussion

### 5.1. Carbon and Nitrogen Abundances-FDU Signature

The carbon and nitrogen abundances (Table 2) reveal signs of the FDU when comparing the results obtained for subgiants and RC stars. As discussed previously, the determination of the nitrogen abundances is only possible for the subgiants and giants. The signature of the FDU is most apparent in the comparison of the ratio ${ }^{12} \mathrm{C} /{ }^{14} \mathrm{~N}$. We find ${ }^{12} \mathrm{C} /{ }^{14} \mathrm{~N}=2.34$ for the subgiants and ${ }^{12} \mathrm{C} /{ }^{14} \mathrm{~N}=1.73$ for the RCs. The decrease of this ratio in the RC stars is indicative of the dredge-up of ${ }^{14} \mathrm{~N}$ that is driven by H burning during the CN cycle. Concerning oxygen, the OH lines become too weak to be useful to measure the oxygen abundances in stars with $T_{\text {eff }}>5000 \mathrm{~K}$. In the case of RCs, however, an oxygen abundance is found with $\langle A(O)\rangle=8.65 \pm 0.01$, consistent with the adopted solar oxygen abundance.

### 5.2. Abundance Variations in M67

The chemical abundances obtained for the eight M67 members studied here, taken at face value, would indicate a measurable abundance spread within M67 for some of the studied elements. However, in all cases, the abundances in the two targets of the same stellar class are found to be quite homogeneous, hinting that when comparing stars across significantly different $T_{\text {eff }}-\log$ regimes, the analyses may be detecting effects other than simple primordial abundance dispersions.

### 5.2.1. Solar-like Stars

The two G dwarfs studied here have atmospheric parameters similar to those of the Sun $\left(T_{\text {eff }}=5724 \mathrm{~K}, \log g=4.48\right.$ for $2 \mathrm{M} 08510076+1153115 ; \quad T_{\text {eff }}=5958 \mathrm{~K}, \quad \log g=4.35$ for $2 \mathrm{M} 08512314+1154049$ ). The chemical similarities between solar-like stars in M67 and the Sun have been discussed previously (Önehag et al. 2014).


Figure 5. Chemical abundances for the two G-type dwarfs relative to Sun ( $[\mathrm{X} / \mathrm{H}]$, with solar results from this study). The solar abundances $([X / H]=0) \pm 0.05$ dex are indicated as dashed lines. This illustrates the close match to a solar abundance pattern (within less than 0.05 dex) of both G dwarfs studied in M67.

The elemental abundances of the two G dwarfs are found to be very consistent with each other, with a mean difference of $+0.01 \pm 0.03$ dex (in the sense of hotter minus cooler star). Figure 5 illustrates the close match to a solar abundance pattern (within less than 0.05 dex) of both solar-like stars in M67. One of the solar-like stars ( $2 \mathrm{M} 08510076+1153115$, or YBP 1514) has an exoplanet detected by Brucalassi et al. (2014, 2016) with a minimum $M_{\text {Jup }}=0.40$. The other solar-like star (2M08512314+1154049, or YBP 1587) was reported in Pasquini et al. (2012) to be an exoplanet host candidate. Our derived stellar parameters and metallicity for 2M08510076 +113115 suggest that this star is a solar twin exhibiting abundance differences relative to the Sun of $\leqslant 0.04$ dex for all elements.

Certain stellar abundance ratios are important for studies of exoplanet properties, such as the C/O ratio, which is one of the factors that affects the ice chemistry in protoplanetary disks (Bond et al. 2010; Teske et al. 2014). Unfortunately, the APOGEE spectra cannot be used to determine precise oxygen abundances in solar-like stars; however, another interesting abundance ratio that affects the interior structure of rocky exoplanets is $\mathrm{Mg} / \mathrm{Si}$ (Delgado Mena et al. 2010; Brewer \& Fischer 2016; and Unterborn \& Panero 2017). We derive $\mathrm{Mg} / \mathrm{Si}=0.93$ for both M67 G-dwarf stars, ${ }^{22}$ compared to our same value of 0.93 obtained for the Sun.

### 5.2.2. Turnoff Stars

We analyzed two stars from the TOP of the HertzsprungRussell (HR) diagram (2M08514122+1154290: $T_{\text {eff }}=6119 \mathrm{~K}$, $\log g=3.91 ; \quad$ and $\quad 2 \mathrm{M} 08505182+1156559: \quad T_{\text {eff }}=6063 \mathrm{~K}$, $\log g=3.87$ ); the abundances of these two stars have a difference $\delta=0.07$ dex for $\mathrm{Mg}, \mathrm{K}$, and Cr , but these are still within the uncertainties in the abundance determinations; see Table 4.

### 5.2.3. Subgiant Stars

The subgiants analyzed in this work are $2 \mathrm{M} 08513540+1157564$ $\left(T_{\text {eff }}=5596 \mathrm{~K}, \log g=3.77\right)$ and $2 \mathrm{M} 08514474+1146460\left(T_{\text {eff }}=\right.$ $5137 \mathrm{~K}, \log g=3.64)$; $2 \mathrm{M} 08513540+1157564$ is located between

[^1]the turnoff of the main sequence and the base of the RGB, while $2 \mathrm{M} 08514474+1146460$ is near the base of the RGB. Their carbon and nitrogen abundances are found to be slightly different in the two stars, and such differences are expected based on stellar evolution models (e.g., Lagarde et al. 2012): the mean difference ( $\pm \mathrm{rms}$ ) between both subgiant stars is $0.02 \pm 0.04 \mathrm{dex}$.

### 5.2.4. Red Clump Stars

The RC stars analyzed here have similar atmospheric parameters: $2 \mathrm{M} 08521856+1144263$ with $T_{\text {eff }}=4842 \mathrm{~K}$, $\log g=2.45$ and $2 \mathrm{M} 08514388+1156425$ with $T_{\text {eff }}=4819 \mathrm{~K}$, $\log g=2.45$. Both stars are members of the RC of M67. As with the other pairs of stellar types isolated here, these M67 RCs share a nearly identical chemistry, with the mean difference in elemental abundances being $+0.02 \pm 0.03$ dex (where the difference is hotter giant-cooler giant). The largest difference is 0.08 dex for ${ }^{14} \mathrm{~N}$, which could be the result of slightly different mixing and mass-loss histories through the FDU and He core flash.

### 5.3. Signatures of Diffusion in M67

Within the pairs of stars of the same stellar classes analyzed here, the chemical compositions are quite homogeneous, while a comparison across the stellar classes may be used to probe the existence and extent of diffusion. Convective mixing predicts (e.g., Lagarde et al. 2012) that red giant photospheres become richer in nitrogen due to internal stellar nucleosynthesis and deep mixing; however, an increase in the abundances of elements such as $\mathrm{Na}, \mathrm{Mg}, \mathrm{Al}, \mathrm{Si}, \mathrm{Ca}, \mathrm{Mn}$, and Fe is not expected in lowmass giants, such as those found in M67. Small increases in red giant abundances relative to those of main-sequence and perhaps subgiant stars might be associated with stellar diffusion operating in the hotter main-sequence stars, while the convective envelopes developing in the atmospheres of evolved stars would tend to erase the diffusion signature. As pointed out in Dotter et al. (2017), the atomic diffusion mechanism operates most effectively in the stars' radiative regions.
Önehag et al. (2014) investigated possible diffusion signatures in M67 by studying a sample of hot, mainsequence stars just below the turnoff ( $T_{\text {eff }} \sim 6130-6200 \mathrm{~K}$ ), turnoff stars ( $T_{\text {eff }} \sim 6150-6215 \mathrm{~K}$ ), and early subgiant stars


Figure 6. Mean abundances for the red clump (red curve), subgiants (blue curve), and turnoff (green curve) minus abundances for the solar twin 2 M 08510076 +1153115 . The red clump and turnoff stars show systematically higher and lower elemental abundances, respectively, compared to the solar twin.
( $T_{\text {eff }} \sim 6040-6110 \mathrm{~K}$ ). Comparing the abundances in the subgiant and turnoff stars, Önehag et al. (2014) found differences in elemental abundances of $\Delta(\mathrm{SG}-\mathrm{TO}) \sim+0.02$ to +0.06 dex; this is in the correct sense predicted by models of diffusion, as the heavy elements that have sunk below the small convection zones in the hotter turnoff stars are mixed back to the observable surface by the deepening convective envelopes in the subgiants. The specific elements studied by Önehag et al. (2014) were C, O, Na, Al, Si, Mg, S, Ca, Ti, Cr, $\mathrm{Mn}, \mathrm{Fe}$, and Ni , and our study includes 10 of these elements in the G dwarfs and subgiants.

Results derived here for the turnoff, subgiant, and RC stars compared to those for the solar twin (2M08510076+1153115) are examined to search for the effects of diffusion in M67. Figure 6 plots the differences in the abundances between the turnoff, subgiants, and RCs (their mean abundances for each class) minus the solar twin. The RC stars display larger abundances in most elements relative to the solar twin, which may reflect the convective erasure of a diffusion signature in the solar twin and even more to the turnoff stars, or the pattern may point to systematic effects in the analysis, as the two types of stars have quite different values of effective temperature and surface gravity (see Section 3).

The abundance differences in Figure 6 between the subgiants and the solar twin are much smaller in comparison with the RCs, with differences typically $<0.05$ dex. The stellar parameters between these stars are much more similar than those for the RC stars. The pattern in the differences is interesting and warrants closer examination. The magnitude of diffusion varies from element to element, and Önehag et al. (2014) also indicated in their Figure 8 that the expected magnitude of diffusion is stronger for the individual abundances of warm main-sequence and turnoff stars in M67. From Önehag et al. (2014), the order of the magnitude of the diffusion differences would be expected to be the largest in Na , $\mathrm{Mg}, \mathrm{Al}, \mathrm{Fe}$, and C , while differences in $\mathrm{Mn}, \mathrm{Cr}$, and Si would be smaller and for Ca and Ti almost nonexistent. Examining the subgiant-solar twin pattern, we note that $\mathrm{Na}, \mathrm{Mg}$, and Al all exhibit relatively strong positive differences, as do $\mathrm{Si}, \mathrm{Ca}, \mathrm{Mn}$, and Fe . Interestingly, Ca and Cr show small negative differences. The derived abundances for the turnoff stars set the lower limits for elemental abundances in M67. For almost
all species, we obtain abundance differences $<0.10$ dex compared to the solar twin, as well as for the other classes. The exceptions were abundances from $\mathrm{K}, \mathrm{Ti}$, and Cr , the elements that showed less change in their abundances as a function of $\log g$. Given that there may be small, systematic effects in the absolute abundance scale as derived for the subgiants, turnoff, and solar twin (a few hundredths of a dex), the overall pattern in Figure 6 may indicate that the signature of diffusion in M67 is stronger in turnoff stars, followed by solarlike stars.

Figure 6 also indicates that we have an increasing overall metallicity from the turnoff stars to the red clump, with the subgiants and the solar-like stars marginally showing similar levels of chemical composition. When computing overall metallicities for each class (using the $[\mathrm{M} / \mathrm{H}]$ as the sum of all elemental abundances), we obtain: solar-like stars $\langle[\mathrm{M} / \mathrm{H}]\rangle=0.03 \pm$ 0.05 , turnoff stars $\langle[\mathrm{M} / \mathrm{H}]\rangle=-0.02 \pm 0.00$, subgiants stars $\langle[\mathrm{M} / \mathrm{H}]\rangle=0.02 \pm 0.07$, and red clump $\langle[\mathrm{M} / \mathrm{H}]\rangle=0.09 \pm$ 0.07 , where the uncertainty here represents the standard deviation ${ }^{23}$ for all elements shown in Figure 6.
Recently, Dotter et al. (2017) published predicted surface abundance changes in stellar models that were computed with the MESA code and included atomic diffusion (Paxton et al. 2011, 2013, 2015; see also Choi et al. 2016 and Dotter 2016). Surface abundances of a number of elements were presented for stars in different evolutionary stages (main sequence, turnoff, and RGB), with metallicities varying from -2.0 to 0.0 and ages from 4.0 to 15.0 Gyr . Their results show that both metallicity and age are important factors in the efficacy of atomic diffusion in altering stellar surface abundances. Of particular relevance for this study of M67 are models that were computed for solar metallicity and an age of 4.0 Gyr. Such models are shown as the solid curves in the different panels of Figure 7, which plot $\Delta[\mathrm{X} / \mathrm{H}]$ versus stellar mass, where $\Delta[\mathrm{X} / \mathrm{H}]=[\mathrm{X} / \mathrm{H}]_{\text {Current }}-[\mathrm{X} / \mathrm{H}]_{\text {Initial }}$ for the elements $\mathrm{Fe}, \mathrm{Mg}, \mathrm{Si}$, and Ca . The filled red symbols represent the elemental abundances derived for the M67 stars, and the pristine Fe abundance for M67 is taken here to be the mean of the K giants, with this value then used as the fiducial point (i.e., $\delta[\mathrm{Fe} / \mathrm{H}]=0.00)$ for the initial cluster value. It should be noted,

[^2]

Figure 7. Atomic diffusion models (solar metallicity and 4.0 Gyr ; Choi et al. 2016; Dotter et al. 2017) for the stellar mass as a function of $\Delta[\mathrm{X} / \mathrm{H}]$ are shown as the gray lines, where $\Delta[\mathrm{X} / \mathrm{H}]$ indicates the current stellar photospheric abundances of the elements $(\mathrm{Fe}, \mathrm{Mg}, \mathrm{Si}$, and Ca$)$ minus the initial cluster composition. The filled red symbols represent the elemental abundances derived for the M67 stars.
however, that the abundances of red clump stars are increased due to the reduction of the H budget caused by H burning, but this effect is estimated as being very small (a difference of less than 0.01 dex; Dotter et al. 2017), and that a better initial cluster abundance would be achieved using M-dwarf abundances. (The M dwarfs in M67 would be too faint $(H \sim 15)$ to be easily observed with APOGEE.)

As can be seen from Figure 7, the values of $\Delta[\mathrm{Fe} / \mathrm{H}]$ as a function of stellar mass predicted by the Dotter et al. (2017) models (with $[\mathrm{Fe} / \mathrm{H}]=0.0$ and age $=4.0 \mathrm{Gyr}$ ) exhibit a behavior that is very similar to that found for the M67 stars as a function of their estimated masses. The behavior of the other elements, $\mathrm{Ca}, \mathrm{Mg}$, and Si , is reminiscent of that of Fe :
there is a pronounced abundance variation as a function of mass in overall agreement with the atomic diffusion models. However, for Mg and Si , the abundances for the red clump stars are $\sim 0.10$ dex higher than the expectations of the models. In addition, the amplitude of the main-sequence turnoff (MSTO) dip is largest for Ca (given the relatively high Ca abundance in the solar-like stars, which is not in good agreement with the models), while for Mg and Si , the MSTO dip is less pronounced and in better agreement with the models.
Elemental abundances derived in the M67 stars covering a range of evolutionary phases suggest that diffusion processes are at work and have been observed in this cluster. Although the number of stars in this initial APOGEE boutique sample is
small, the comparison of spectroscopically derived Fe abundances with those predicted by stellar models that include atomic diffusion is very promising, and this pilot study demonstrates what can be accomplished with the APOGEE spectra. D. Souto et al. (2018, in preparation) will present abundances for a much larger sample of M67 members based on the stellar parameters obtained from photometry and isochrones and chemical abundances derived automatically from the APOGEE spectra; these results will also be compared with the DR14 abundances.

## 6. Summary

This paper presents detailed chemical abundances for the elements C, N, O, Na, Mg, Al, Si, K, Ca, Ti, V, Cr, Mn, and Fe in eight stellar members of the open cluster M67. The sample stars have different stellar masses and are in different evolutionary stages: two G dwarfs, two turnoff stars, two G subgiants, and two K giants (RC). Abundances were derived via a manual "boutique" detailed spectroscopic line-by-line abundance analysis using APOGEE high-resolution NIR spectra.

The derived abundances were investigated as a function of stellar evolutionary state, with homogeneous abundances found for each pair of stars in the same phase of evolution. Significant changes in abundance ( $\sim 0.05-0.20$ dex) were found across different evolutionary states, with most of the studied elements (except $\mathrm{K}, \mathrm{Ti}$, and Cr ) having their lowest values in the turnoff stars, while the RC stars tended to exhibit the largest abundances. For iron, in particular, we obtain $\langle A(\mathrm{Fe})\rangle=$ $7.48 \pm 0.05$ for solar-like stars, $\langle A(\mathrm{Fe})\rangle=7.39 \pm 0.05$ for turnoff stars, $\langle A(\mathrm{Fe})\rangle=7.46 \pm 0.08$ for subgiant stars, and $\langle A(\mathrm{Fe})\rangle=7.50 \pm 0.07$ for RC stars.

We have conducted a comparison of the derived abundances of $\mathrm{Mg}, \mathrm{Si}, \mathrm{Ca}$, and Fe as a function of stellar mass with atomic diffusion models in the literature (Choi et al. 2016; Dotter 2016; Dotter et al. 2017). We find that the atomic diffusion models explain reasonably well most of the observed abundance trends, suggesting that the signature of diffusion has been detected in M67 stars.

Published departures from LTE-derived abundances from Na I, Mg I, and Si I lines in the APOGEE spectral region were examined and are expected to be significantly smaller than the changes in abundance found in the M67 stars: non-LTE is unlikely to explain the observed abundance spreads that correlate with evolutionary state, but 3D calculations are still needed. Diffusion effects will be investigated further using the entire set of APOGEE spectra and ASPCAP results for M67 members in a follow-up paper.

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Facility: Sloan.
Software: Turbospectrum (Plez 2012), (Alvarez \& Plez 1998), Numpy (Van Der Walt et al. 2011), Matplotlib (Hunter 2007).

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[^0]:    21 marcs.astro.uu.se

[^1]:    $\overline{22} \mathrm{Mg} / \mathrm{Si}=N(\mathrm{Mg}) / N(\mathrm{Si})=10^{\log N(\mathrm{Mg})} / 10^{\log } N(\mathrm{Si})$.

[^2]:    ${ }^{23}$ The uncertainty in the mean is a factor $1 / \operatorname{sqrt}(2)$ smaller.

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