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# Synthesis, characterization and electrical properties of quaternary selenodiphosphates: $AMP_2SE_6$ with A = Cu, Ag and M = Bi, Sb

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#### Abstract

The new quaternary selenophosphate phases  $AMP_2Se_6$  (A = Cu, Ag and M = Bi, Sb) were synthesized by ceramic methods at 1023 K. These phases were characterized by powder X-ray diffraction (XRD), Fourier transform infrared spectroscopy (FTIR), energy dispersive X-ray analysis (EDX) and a.c. and d.c. electrical conductivity measurements. The phases all show values of electrical conductivity,  $\sigma$ , of about  $10^{-4} \Omega^{-1} \text{ cm}^{-1}$  at 303 K and photoconductive effect. The conductivity is nearly five orders of magnitude larger than that of related phases.  $\bigcirc$  2003 Elsevier Science Ltd. All rights reserved.

Keywords: D. Electrical properties; A. Layered compounds

#### 1. Introduction

The study of compounds in the family of  $M_2P_2Q_6$  (M = Mn, Fe, Co, Ni, Cd, V, Mg, Pb, Hg, Sn; Q = S, Se), has been stimulated because of their highly anisotropic physical properties, useful for possible applications in electrooptical devices. The first layered selenodiphosphate synthesized, Fe<sub>2</sub>P<sub>2</sub>Se<sub>6</sub>, was reported by Klingen et al. in 1973 [1]. In the layered structure of this material, the transition metal cation Fe<sup>2+</sup> and the P–P pairs, occupy the octahedral holes between the planes of selenium atoms. The substitution of the cations M<sup>2+</sup>, by M<sup>1+</sup> (Ag, Cu) and M<sup>3+</sup> (such as Cr, In, V, Ga, Sc) results in thermodynamically stable quaternary chalcophosphates with structures of the CdCl<sub>2</sub> type [2].

The selenodiphosphate compounds containing Bi and Sb, have been synthesized by solid state reaction of the elemental powders at low temperature. The phases  $M_4(P_2Se_6)_3$  (M = Bi, Sb) exhibit structures related to Pb<sub>2</sub>P<sub>2</sub>Se<sub>6</sub> [3], while the phase  $\beta$ -Bi<sub>4</sub>(P<sub>2</sub>Se<sub>6</sub>)<sub>3</sub> exhibits a three-dimensional (3D)

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structure, with groups  $BiSe_8$  and  $P_2Se_6^{4-}$  [4]. The quaternary selenodiphosphates  $KMP_2Se_6$  and  $Cs_8M_4(P_2Se_6)_5$  (M = Bi, Sb) have been synthesized using low temperature methods [5]. The phases  $KMP_2Se_6$  are isostructural and exhibit two-dimensional (2D) structures. The phases  $Cs_8M_4(P_2Se_6)_5$  present tunnels interconnected by  $M \cdots M$  interactions in a 2D structure. An important characteristic of the Bi and Sb atoms, in solid phases, is the stereochemical localization of the  $ns^2$  electron pairs, which produces important variations in crystal structure, physical properties and electronic structure of the resulting materials.

In this work, we report the synthesis, characterization and electrical properties of the new phases  $AMP_2Se_6$  (A = Cu, Ag; M = Bi, Sb), which are structurally related to  $Fe_2P_2S_6$ . As far as we know, up to now, these phases are the first quaternary compounds of Bi and Sb that contain a transition metal counterion (A<sup>1+</sup>).

## 2. Experimental section

#### 2.1. Synthesis

The phases AMP<sub>2</sub>Se<sub>6</sub> (A = Cu, Ag; M = Bi, Sb) were synthesized by two sets of experimental conditions. In both, the reaction mixtures were sealed in evacuated quartz ampoules, and heated at 1023 K for 1 week. In the first method, the compounds were obtained by direct combination of the pure elements in stoichiometric proportions, with an excess of selenium and phosphorus (1% mass) to avoid the formation of impurity phases (Cu<sub>3</sub>PSe<sub>4</sub>, Ag<sub>4</sub>P<sub>2</sub>Se<sub>6</sub> and M<sub>2</sub>Se<sub>3</sub>). Homogeneous and well-crystallized materials were obtained after grinding and reheating at 1023 K for 1 week. SEM-EDX analyses carried out on the samples confirmed their purity, homogeneity and stoichiometry (Table 1). In the second procedure, the synthesis proceeded by direct combination of the pure elements with an excess of P<sub>2</sub>Se<sub>5</sub> and Se, so that all of the metal reacted to give well-crystallized products. A, M, P<sub>2</sub>Se<sub>5</sub>, and Se were mixed in molar ratios 1:1:1.5:2.0. After the reaction was completed, the reacted matter was slowly cooled to room temperature at the rate of 6 K/h. The reaction mixture was washed with DMF/ ethylenediamine (3:1) to remove the P<sub>y</sub>Se<sub>z</sub> and Se. The product was washed with anhydrous ether and consisted of black laminar crystals. EDX analysis carried out on the resulting crystals indicated an average composition of Cu<sub>1.0</sub>Bi<sub>1.0</sub>P<sub>1.8</sub>Se<sub>6.1</sub>.

#### 2.2. Characterization

XRD data were collected at room temperature on a Siemens D 5000 powder diffractometer, with Cu K $\alpha$  radiation in the range 5 < 2 $\theta$  < 80. The lattice parameters were calculated from least-square fits

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Cu	Ag	Bi	Sb	Р	Se	Proposed formula
7.8	-	25.7	_	7.3	59.2	CuBiP <sub>2</sub> Se <sub>6</sub>
-	13.2	24.2	_	7.2	55.4	AgBiP <sub>2</sub> Se <sub>6</sub>
10.0	_	_	16.5	8.2	65.3	CuSbP <sub>2</sub> Se <sub>6</sub>
-	14.3	_	16.1	7.8	61.8	AgSbP <sub>2</sub> Se <sub>6</sub>

Table 1 Analytical data (% w/w) for the AMP<sub>2</sub>Se<sub>6</sub> phases

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using the Powder Diffraction Package (PDP) and Powder Pattern Lattice Parameter (PPLP) refinement routine of the NRCVAX program. For electron microscopy, analyses were performed on a JEOL 6400 scanning electron microscopy (SEM) equipped with Oxford Link Isis energy dispersive X-ray (EDX) detector. The average A:M:P:Se elemental ratio of all samples was 1:1:2:6.

Differential Thermal Analysis (DTA) and Thermogravimetric Analysis (TG) were performed on a Rheometric Scientific STA 1500H/625 Thermal Analysis System. The DTA/TG curves were run simultaneously on each sample from room temperature to 1273 K, in flowing nitrogen at a heating rate of 10 K/min. Infrared spectra were recorded on a BRUKER VECTOR 22 with Fourier transform SJ-IR spectrometer in the spectral range  $250-600 \text{ cm}^{-1}$ . The spectra of the powders were recorded in polyethylene pellets. Magnetic susceptibility ( $\gamma$ ) measurements were performed on polycrystalline samples between 5 K and room temperature under a constant field strength of 1 kG, using a SHE VTS-906 SQUID detection system. The electrical conductivity was measured by a.c. and d.c. methods. The two opposite flat surfaces of the samples were sputtered with gold and sandwiched between the platinum electrodes of the conductivity cell. Impedance measurements were carried out with a Solartron SI-1260 impedance gain-phase analyzer in the frequency range 0.1–10 MHz with a signal level between 25 mV and 1 V, to probe the charge transport mechanism. Direct current measurements were performed with a Keithley 237 source-meter, to ascertain the existence of slow relaxation mechanisms that may not be accessible through the impedance measurements. Photoconductivity measurements were carried out in the temperature range 296-473 K, and pressure of approximately  $10^{-3}$  atmosphere. The samples, pressed into cylindrical disks, were mounted in a metallic holder, where illumination could be achieved through a transparent window. The light power density was 50 mW/cm<sup>2</sup> from a Doolan Jenner Fiberlite A3200 halogen lamp. The current was recorded with a programmable Keithley 617 electrometer with silver paint contacts and planar configuration of the electrodes. The I-Vcurves were verified to be linear, supporting the Ohmic character of the contacts.

## 3. Results and discussion

The indexing of the X-ray powder diffractograms shows that the phases AMP<sub>2</sub>Se<sub>6</sub> belong to the space group *C2/m* and their structures are related to that of Fe<sub>2</sub>P<sub>2</sub>S<sub>6</sub>. The phases AMP<sub>2</sub>Se<sub>6</sub> present a 2D structure and they can be regarded as a substitution of Fe<sup>2+</sup>, in the Fe<sub>2</sub>P<sub>2</sub>S<sub>6</sub> layer, by the cations A<sup>1+</sup> and M<sup>3+</sup>. The observed and calculated interplanar spacing ( $d_{hkl}$ ) and the unit cell constants are shown in Tables 2 and 3, respectively. The 0 0 1 interplanar *d*-spacing shows that the phases correspond to unintercalated layer compounds. The infrared spectrum shows a strong absorption at 428 cm<sup>-1</sup> in CuBiP<sub>2</sub>Se<sub>6</sub>, 430 cm<sup>-1</sup> in AgBiP<sub>2</sub>Se<sub>6</sub>, 432 cm<sup>-1</sup> in CuSbP<sub>2</sub>Se<sub>6</sub> and 429 cm<sup>-1</sup> AgSbP<sub>2</sub>Se<sub>6</sub>. This vibration can be assigned to the asymmetric stretching ( $\nu$ PSe<sub>3</sub>), indicating the presence of P<sub>2</sub>Se<sub>6</sub><sup>4-</sup> units in the structure. Both IR and XRD results indicate that the metals A<sup>1+</sup> and M<sup>3+</sup> were incorporated, rather than intercalated, into the layer of the phases AMP<sub>2</sub>Se<sub>6</sub>. Quaternary phases of bismuth and antimony reported in the literature also possess layer structures. The phase KMP<sub>2</sub>Se<sub>6</sub> (M = Sb, Bi) contains K<sup>+</sup> intercalated in the interlayer space of lamellar MP<sub>2</sub>Se<sub>6</sub><sup>1-</sup>, where Cs<sub>8</sub>M<sub>4</sub>(P<sub>2</sub>Se<sub>6</sub>)<sub>5</sub> (M = Bi, Sb) present a staircase layered structure with tunnels along the *b* direction, Cs<sup>1+</sup> is found in the tunnels and within the layer.

DTA/TG showed a slight weight loss of approximately 2% between 523 and 673 K, which may be assigned to partial selenium loss, and a decomposition peak at 803 K with a weight loss of 23.8%. The XRD of the residues indicated the presence of  $M_2Se_3$  and  $A_2Se$ .

Table 2					
Observed	and	calculated	Powder	Diffraction	data

CuBiP <sub>2</sub> Se <sub>6</sub>				AgBiP <sub>2</sub> Se <sub>6</sub>					
$2\theta$	Int	$d_{\mathrm{cal}}$	$d_{\rm obs}$	h k l	$2\theta$	Int	$d_{\mathrm{cal}}$	$d_{\rm obs}$	h k l
13.58	100	6.51	6.51	001	13.50	100	6.56	6.55	001
17.29	3	5.13	5.13	110	16.80	1	5.28	5.27	020
18.83	6	4.708	4.708	$-1\ 1\ 1$	17.87	1	4.96	4.96	110
20.86	1	4.256	4.256	021	19.12	1	4.641	4.637	-1111
27.12	58	3.285	3.285	002	20.84	3	4.262	4.259	021
29.65	4	3.011	3.011	-201	24.39	1	3.649	3.647	111
30.67	26	2.913	2.913	-131	27.08	42	3.292	3.289	002
31.77	2	2.814	2.814	022	27.42	11	3.253	3.250	-112
32.42	2	2.759	2.759	040	29.82	9	2.996	2.993	-201
38.82	16	2.138	2.138	-222	30.18	19	2.961	2.959	200
39.55	1	2.277	2.277	-113	32.52	1	2.755	2.753	022
40.62	3	2.219	2.219	221	33.29	1	2.692	2.689	040
41.05	23	2.197	2.197	003	38.48	10	2.340	2.337	-222
48.26	5	1.884	1.884	-242	39.67	11	2.272	2.270	220
49.93	9	1.825	1.825	222	41.05	19	2.199	2.197	003
51.00	3	1.789	1.789	-331	44.69	1	2.028	2.026	023
55.60	41	1.652	1.651	004	45.41	2	1.997	1.995	240
57.45	1	1.603	1.603	-224	49.71	7	1.834	1.832	222
58.81	2	1.569	1.569	-2.61	55.67	23	1.651	1.648	004
60.25	- 1	1.5349	1.5347	-402	60.63	1	1.5272	1.5250	260
61.22	1	1.5129	1.5129	223	62.80	2	1.4797	1.4782	$\frac{1}{400}$
65.69	2	1.4203	1.4202	171	67.42	-	1.3890	1.3887	332
71.07	- 1	1 3253	1 3250	-2.2.5	70.91	2	1 3290	1 3285	-442
76.53	3	1.2439	1.2439	082	76.14	-	1.2502	1.2516	2.2.4
77.74	6	1.2274	1.2274	-2.81	77.70	4	1.2289	1.2288	-512
,,.,.	0	1.2271	1.227	201	79.07	2	1 2111	1 2107	-335
	_					-		112107	000
CuSbP <sub>2</sub>	Se <sub>6</sub>				AgSbP <sub>2</sub> S	e <sub>6</sub>			
13.63	50	6.49	6.49	001	13.36	100	6.62	6.62	001
17.21	1	5.15	5.15	110	16.92	3	5.23	5.23	110
19.31	2	4.594	4.594	-111	21.53	2	4.125	4.124	021
19.62	1	4.521	4.520	021	26.85	64	3.318	3.317	002
24.17	1	3.679	3.679	111	30.31	13	2.946	2.946	200
27.17	100	3.280	3.279	002	31.23	3	2.861	2.861	022
27.76	5	3.212	3.211	-112	34.17	4	2.622	2.621	-202
29.68	2	3.008	3.008	-201	35.71	2	2.512	2.512	112
30.89	23	2.892	2.892	-131	37.21	3	2.414	2.414	-222
31.86	1	2.806	2.805	022	40.17	17	2.243	2.214	-113
34.44	2	2.602	2.601	-202	41.40	3	2.179	2.179	003
39.02	12	2.306	2.306	-222	43.73	4	2.068	2.068	023
39.85	3	2.260	2.263	-113	47.83	6	1.900	1.901	-133
41.07	8	2.196	2.196	003	49.50	8	1.836	1.836	222
44.27	1	2.045	2.044	023	51.55	3	1.772	1.771	330
46.00	1	1.972	1.972	-151	55.24	22	1.661	1.661	-204
48.73	6	1.867	1.867	113	57.65	3	1.598	1.597	-224
50.11	11	1.819	1.818	222	62.59	2	1.4828	1.4828	400
55.63	28	1.651	1.650	004	63.97	2	1.4541	1.4541	-352

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CuBiP <sub>2</sub> S	e <sub>6</sub>				AgBiP <sub>2</sub> Se <sub>6</sub>				
20	Int	$d_{\rm cal}$	$d_{\rm obs}$	h k l	$2\theta$	Int	$d_{\rm cal}$	$d_{\rm obs}$	h k l
63.19	2	1.4703	1.4702	-262	69.16	2	1.3572	1.3571	351
65.34	2	1.4271	1.4272	420	75.97	1	1.2515	1.2514	224
76.75	2	1.2408	1.2407	082	77.39	3	1.2321	1.2321	-354
77.86	7	1.2259	1.2258	-281					
79.49	2	1.2047	1.2047	-372					

Table 2 (Continued)

Table 3 Lattice parameters

Phase	a (Å)	<i>b</i> (Å)	c (Å)	β (°)	$V(\text{\AA}^3)$
CuBiPoSec	6 19 (2)	10.73 (2)	6 94 (2)	107.46 (6)	439 37
AgBiP <sub>2</sub> Se <sub>6</sub>	6.20 (2)	10.71 (3)	6.90 (2)	107.20 (7)	437.55
$CuSbP_2Se_6$	6.16 (2)	10.71 (4)	6.89 (2)	106.48 (6)	436.13
AgSbP <sub>2</sub> Se <sub>6</sub>	6.20 (2)	10.73 (3)	6.90 (2)	107.34 (9)	437.88

Magnetic measurements carried out on  $CuSbP_2Se_6$  showed temperature-independent paramagnetism. The calculated magnetic moment is 0.19 BM, which could be related to non-stoichiometry in the Cu sublattice.

The compounds  $AMP_2Se_6$  are semiconductors with electrical conductivity,  $\sigma$ , of about  $10^{-4} \Omega^{-1} \text{ cm}^{-1}$  at 303 K (Table 4). These values are much larger than those of the related compounds  $Cs_8M_4(P_2Se_6)_5$  (M = Sb, Bi). Thus,  $Cs_8Sb_4(P_2Se_6)_5$  and  $Cs_8Bi_4(P_2Se_6)_5$  are insulatings, the bismuth compound shows electrical conductivity of  $\sigma = 10^{-9} \Omega^{-1} \text{ cm}^{-1}$  at room temperature [5b]. From the above results, it appears that the presence of the cations  $Ag^+$  and  $Cu^+$  in  $AMP_2Se_6$  is responsible for this large increase in electrical conductivity.

The a.c. impedance measurements give typical depressed semicircular plots (Nyquist diagrams), suggesting the presence of wide distributions of relaxation times. Fig. 1 shows representative plots of CuBiP<sub>2</sub>Se<sub>6</sub> and AgBiP<sub>2</sub>Se<sub>6</sub> (parts a and b, respectively) measured at room temperature. By fitting discrete elements models [6,7] it is possible to identify the resistance and effective capacitance corresponding to the single arc, resulting in the values reported in Table 5. The values of the effective capacitance suggest that the electrical conductivity in these AMP<sub>2</sub>Se<sub>6</sub> compounds is controlled by the intergranular transport. The Nyquist plot resulting from AgBiP<sub>2</sub>Se<sub>6</sub> has been singled out in Fig. 1b, as it

Table 4 Dark conductivity ( $\sigma_d$ ) and photoconductivity ( $\sigma_{ph}$ ) at 303 K

Phase	$\sigma_{\rm d}~(\Omega^{-1}~{ m cm}^{-1})$	$\sigma_{\rm ph}~(\Omega^{-1}~{\rm cm}^{-1})$	$\sigma_{ m ph}/\sigma_{ m d}$
CuBiP <sub>2</sub> Se <sub>6</sub>	$1.46 \times 10^{-4}$	$1.94 \times 10^{-4}$	1.33
AgBiP <sub>2</sub> Se <sub>6</sub>	$6.34 \times 10^{-4}$	$6.59 \times 10^{-4}$	1.04
CuSbP <sub>2</sub> Se <sub>6</sub>	$4.56 \times 10^{-5}$	$5.96 \times 10^{-5}$	1.20
AgSbP <sub>2</sub> Se <sub>6</sub>	$1.45 \times 10^{-4}$	$1.65 \times 10^{-4}$	1.14



Fig. 1. Impedance complex plane plot of (a) CuBiP<sub>2</sub>Se<sub>6</sub>, (b) AgBiP<sub>2</sub>Se<sub>6</sub>.

is the only one to deviate noticeably from the single depressed arc going through the origin behaviour. Indeed, the arc, which starts out at the maximum frequency available from the experimental system, is displaced from the origin. The series combination of a resistance  $(R_{\rm gi})$  with the parallel combination of a resistance  $(R_{\rm gb})$  and a constant phase element  $(C_{\rm gb})$  fits the data. The resistance  $R_{\rm gi}$ , which causes the

Phase	Temperature (K)	$R_{\rm gb} (\Omega)$	$C_{\rm gb}$ (F)
CuBiP <sub>2</sub> Se <sub>6</sub>	293	3040	$3.01 \times 10^{-10}$
AgBiP <sub>2</sub> Se <sub>6</sub>	288	$R_{\rm gi}$ : 864, $R_{\rm gb}$ : 1593	$7.10 \times 10^{-11}$
CuSbP <sub>2</sub> Se <sub>6</sub>	295	17890	$1.93 \times 10^{-11}$
AgSbP <sub>2</sub> Se <sub>6</sub>	297	3125	$3.20 \times 10^{-11}$

 Table 5

 Resistance and effective capacitance calculated from the Nyquist curves

displacement, is 864  $\Omega$  and the second, defining the arc diameter, is 729  $\Omega$ . The value of the effective capacitance, 71 pF, in the latter element, suggests control by intergranular transport, as in the other samples, however, the displacement represented by  $R_{gi}$  suggests that an arc due to intragrain transport may be active at higher frequency than registered. Both these values of the resistance are reported in Table 5, leading to an effective d.c. resistance of 1593  $\Omega$ .

The d.c. electrical conductivity measurements show the characteristic behavior of photoconductive materials [8]. Fig. 2 shows the current as a function of time, where a monotonic rise to a steady value can be observed upon illumination, followed by quite a slow decay of the photoconductivity upon removal of the illumination source. A persistent photoconductivity effect is observed in AgBiP<sub>2</sub>Se<sub>6</sub> after turning off the illumination. This type of photoconductivity behavior has been reported in chalcogenide glasses and intercalated chalcogenides [9]. In all of the cases where thermally activated behavior is observed, the experimental data correlate with the expression:

$$\sigma = \sigma_0 \exp\left(\frac{-E_a}{kT}\right)$$

Table 6

where  $\sigma_0$  is a constant,  $E_a$  is the activation energy, k is Boltzmann's constant and T is the absolute temperature.

The temperature dependence of the d.c. conductivity presents approximately Arrhenius behavior in the temperature range of the measurements, as indicated in Fig. 3. In all four phases, the dark conductivity increases as the temperature rises. This behavior is consistent with the thermal activation of trapped carriers in a wide-gap semiconductor material. The corresponding activation energies are collected in Table 6. Both dark conductivity  $\sigma_d$  and photoconductivity  $\sigma_{ph}$  are thermally activated, in the whole temperature range investigated. In the phases AMP<sub>2</sub>Se<sub>6</sub> studied, those with A = Cu show greater photosensitivity ( $\sigma_{ph}/\sigma_d$ ) than those with A = Ag (Table 4).

Activation a	noral	and tam	noroturo	rongo	of dorl	, aanduativit	and.	nhotocondu	ativity
Activation	nergy a	and tem	perature	Tange (	uan uan	conductivit	y anu	photocolidu	icuvity

Phase	Temperature range (K)	$E_{\rm a}~({\rm eV})$
Dark conductivity		
AgBiP <sub>2</sub> Se <sub>6</sub>	303–375	0.06
$AgSbP_2Se_6$	328–373	0.29
$CuSbP_2Se_6$	303-403	0.15
CuBiP <sub>2</sub> Se <sub>6</sub>	303–463	0.14
Photoconductivity		
CuBiP <sub>2</sub> Se <sub>6</sub>	303–463	0.10



Fig. 2. Direct current photoconductivity as a function of time: (a) CuBiP<sub>2</sub>Se<sub>6</sub>, (b) AgBiP<sub>2</sub>Se<sub>6</sub>.

For interpreting the photoconductivity in the phases  $AMP_2Se_6$ , we have used an extended Huckel (EH) calculation scheme based on the crystal structure reported for  $CuCrP_2Se_6$ , isostructural with  $Fe_2P_2S_6$  [2a]. In these schemes, the highest energy valence bands are derived from the (P–P) p bonding orbitals, from (Se) p nonbonding orbitals and from (P–Se) p hybridized orbitals. The lowest energy conduction bands are generated by the (Cu) s bonding orbitals, the (P–P) p antibonding and (Se) p bonding states. The (Bi) s orbitals are below the main valence band. The exposure of the sample to light produces holes in the (Se) p valence nonbonding band and electrons in the (Cu) s conduction bonding band. The free centers in the (P–P) p valence bonding band act as recombination centers.



Fig. 3. Direct current dark conductivity and photoconductivity as a function of temperature: (a) photoconductivity of  $CuBiP_2Se_6$ , (b) dark conductivity of  $CuBiP_2Se_6$ , (c) dark conductivity of  $AgBiP_2Se_6$ , (d) dark conductivity of  $AgSbP_2Se_6$  and (e) dark conductivity of  $CuSbP_2Se_6$ .

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